

# MOS MODEL 9

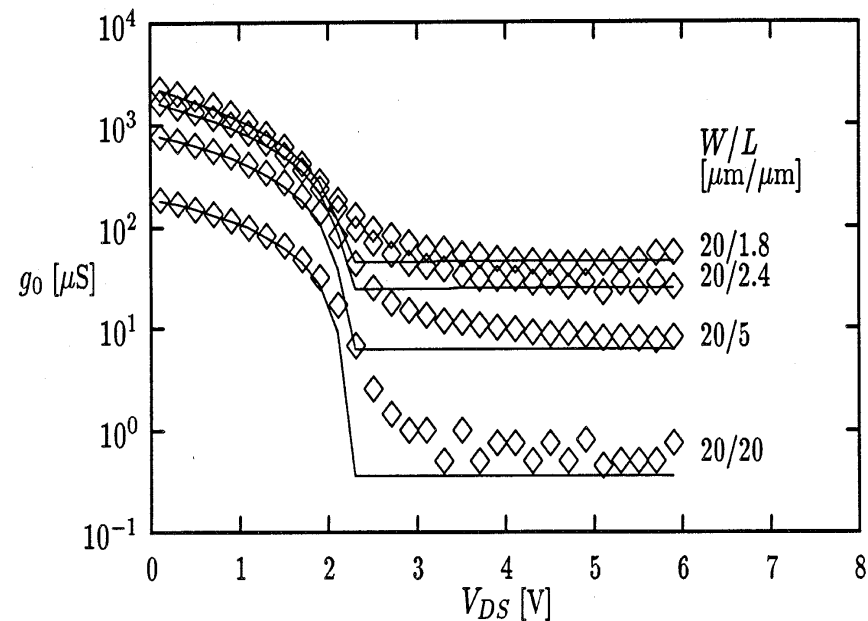
## from Physics to Equations

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- introduction
- linear region
- subthreshold region
- saturation region
- output conductance
- avalanche or substrate current
- summary

- MOS MODEL 9 has been introduced in 1990 on C3DM
- before 1990 MOS MODEL 7 was the most advanced compact model

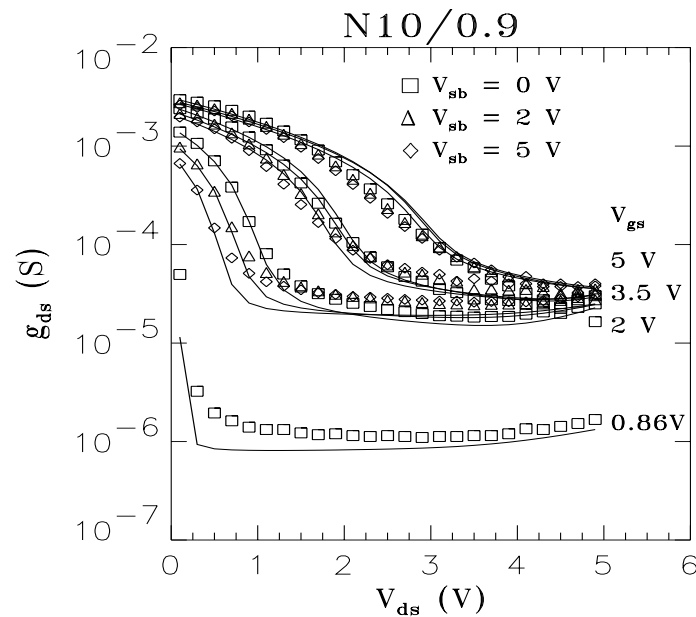


MOS MODEL 7  
output conductance

- MOS MODEL 7 is only suited for digital applications

## MOS MODEL 9

- has been developed for analog applications

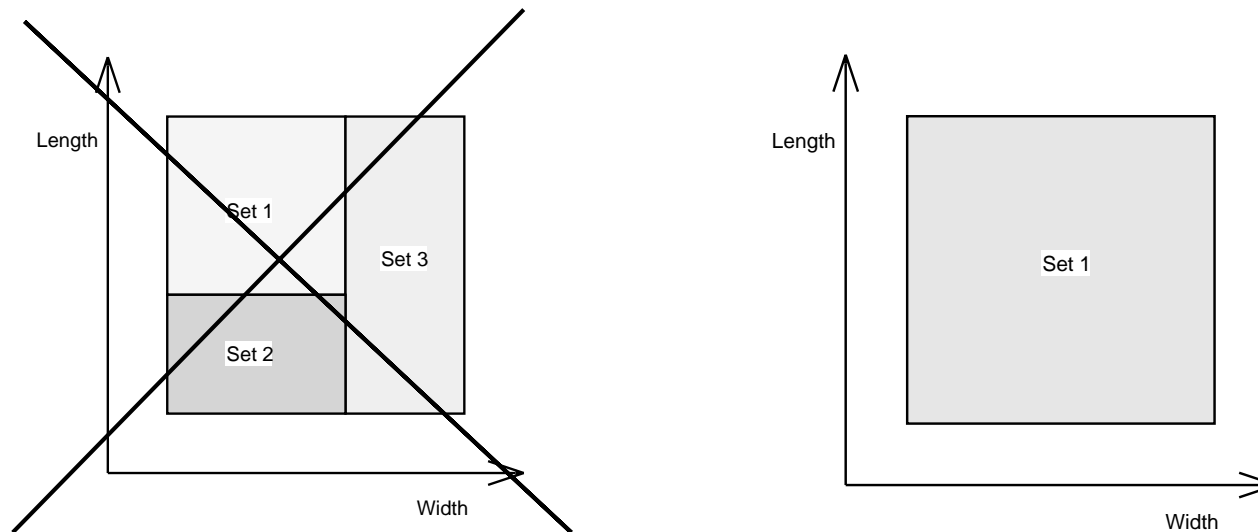


MOS MODEL 9  
output conductance

- has continuous first and second order derivatives
- is suited for both digital and analog applications

## MOS MODEL 9

- describes whole geometry range of process with **one** parameter set



- includes temperature scaling
- has no additional parameters for the charge model

## MOS MODEL 9

- relation between physics and equations
  - parameters
  - scaling rules
- no exact, detailed derivations
- trends in specific parameters

- drift current  $I_{ds}$  from continuity equation

$$I_{ds} = -\mu W_{\text{eff}} Q_i \frac{\partial \phi}{\partial x}$$

- integration over  $x$

$$\int_0^{L_{\text{eff}}} I_{ds} dx = L_{\text{eff}} I_{ds} = -\mu W_{\text{eff}} \int_0^{V_d} Q_i d\phi$$

- intermediate result

$$I_{ds} = -\mu \frac{W_{\text{eff}}}{L_{\text{eff}}} \int_0^{V_d} Q_i d\phi$$

- $\mu \Rightarrow$  carrier mobility;  $Q_i \Rightarrow$  inversion charge per unit surface

- intermediate result

$$I_{ds} = -\mu \frac{W_{\text{eff}}}{L_{\text{eff}}} \int_0^{V_d} Q_i d\phi$$

- inversion charge  $Q_i$  from Poisson equation

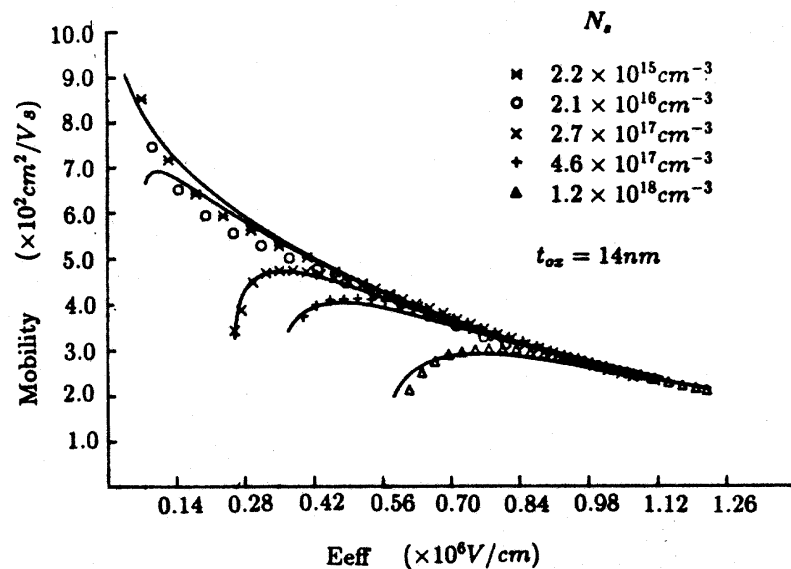
$$Q_i \approx -C_{\text{ox}} (V_{\text{gs}} - V_{\text{T0}} - \phi)$$

- substitution and integration

$$I_{ds} = \mu C_{\text{ox}} \frac{W_{\text{eff}}}{L_{\text{eff}}} \left[ (V_{\text{gs}} - V_{\text{T0}}) V_d - \frac{1}{2} V_d^2 \right]$$

- $C_{\text{ox}} \Rightarrow$  oxide capacitance per unit surface  
 $V_{\text{T0}} \Rightarrow$  MM9 threshold voltage

- carrier mobility depends on electric field perpendicular to surface



carrier mobility vs.  $E_{\perp}$

- mobility is a function of gate-source voltage

$$\mu = \frac{\mu_0}{1 + \theta_1 (V_{\text{gs}} - V_{\text{T0}})}$$

- $\theta_1 \Rightarrow$  MM9 mobility reduction parameter

- MM9 expression for  $I_{ds}$

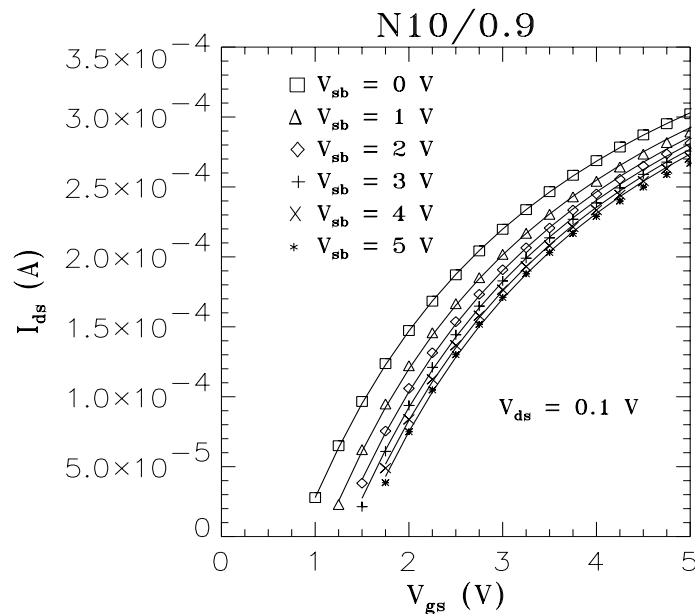
$$I_{ds} = \beta \frac{(V_{gs} - V_{TO}) V_{ds} - \left(\frac{1 + \delta_1}{2}\right) V_{ds}^2}{1 + \theta_1 (V_{gs} - V_{TO})}$$

- $\beta$  is the MM9 gain factor

$$\beta = \mu_o C_{ox} \frac{W_{eff}}{L_{eff}}$$

- charge in depletion layer increases with  $V_{ds}$

$$\delta_1 \approx 0.3 \frac{k}{\sqrt{V_{sb} + \phi_b}}$$



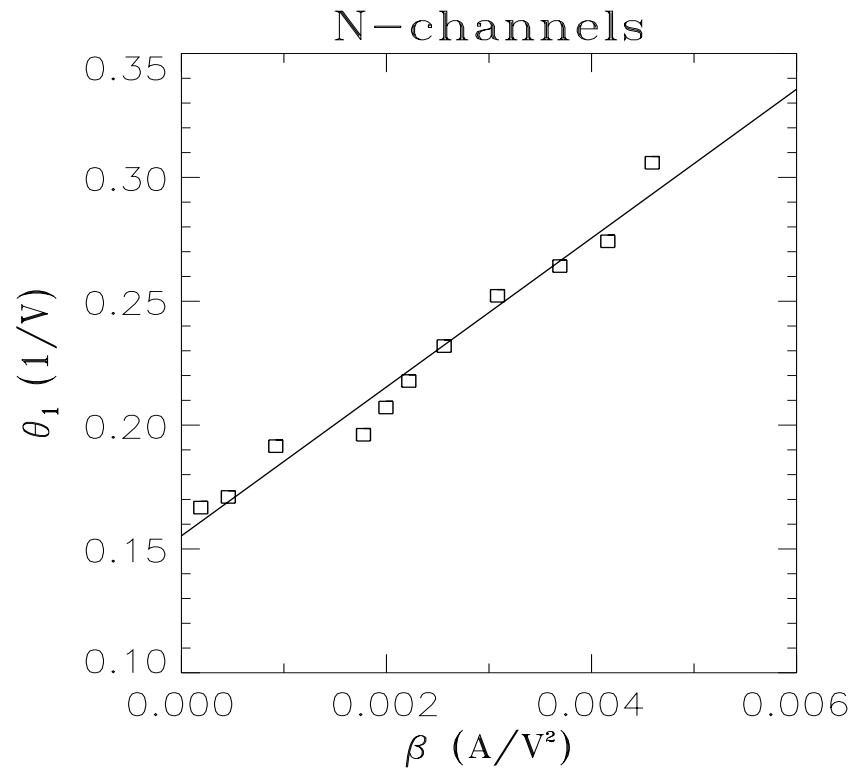
MOS MODEL 9  
linear region

- MM9 expression for  $I_{ds}$

$$I_{ds} = \beta \frac{(V_{gs} - V_{TO}) V_{ds} - \left(\frac{1 + \delta_1}{2}\right) V_{ds}^2}{1 + \theta_1 (V_{gs} - V_{TO})}$$

- source and drain series resistances are included in  $\theta_1$

$$\theta_1 = \theta_{10} + \beta (R_s + R_d)$$

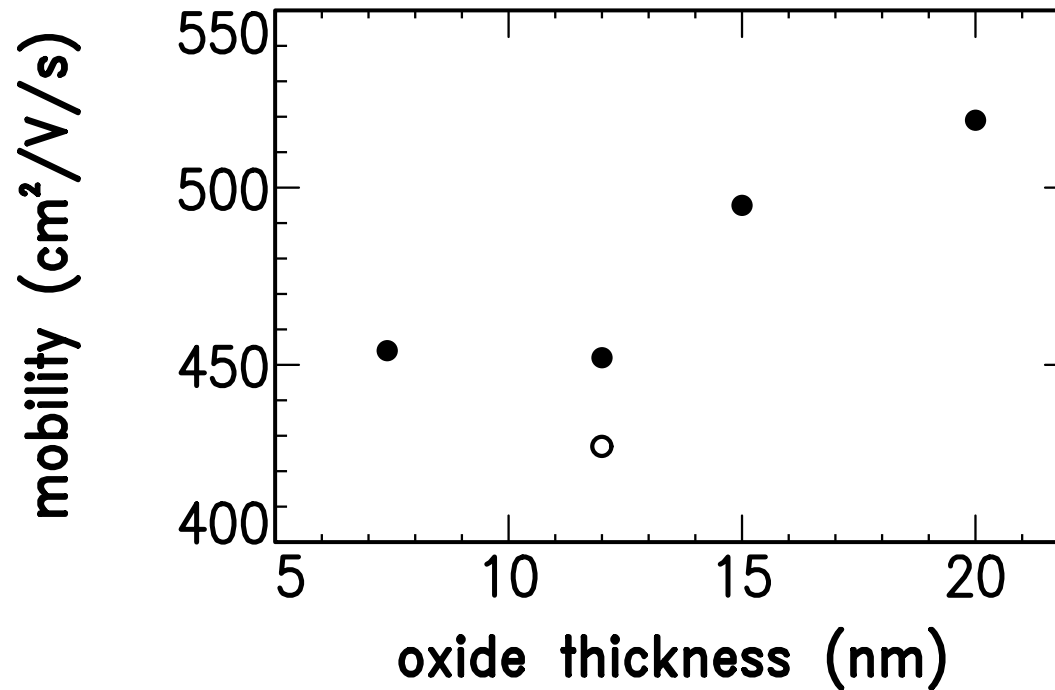


MOS MODEL 9

$\theta_1$  vs.  $\beta$

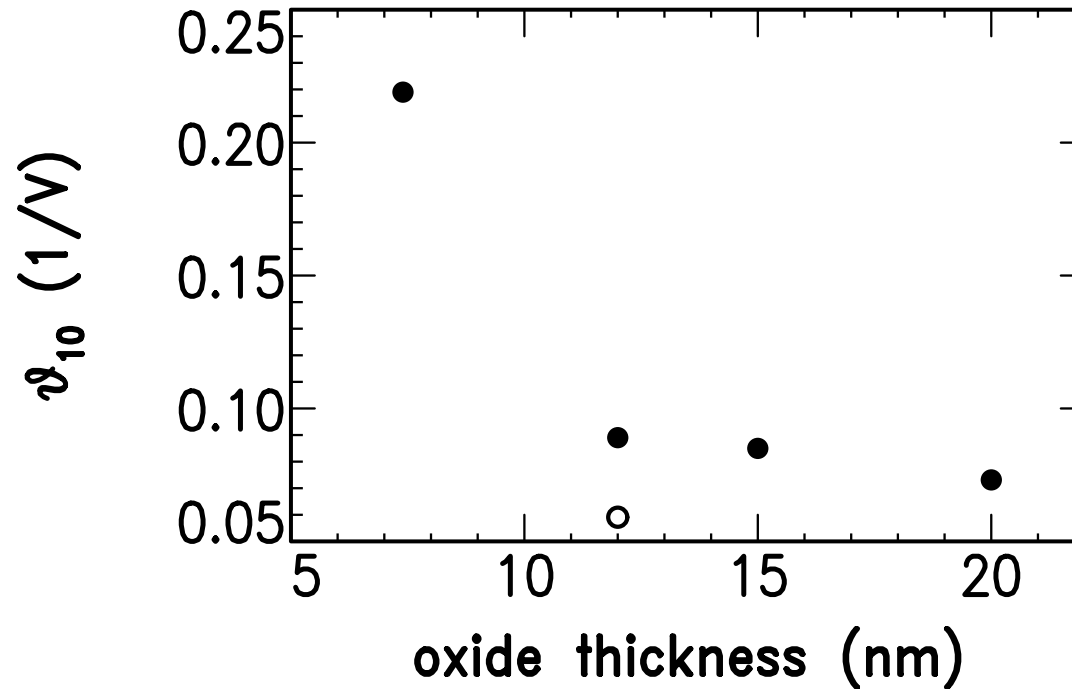
- source and drain series resistances are included in  $\theta_1$

$$\theta_1 = \theta_{10} + \beta (R_s + R_d)$$

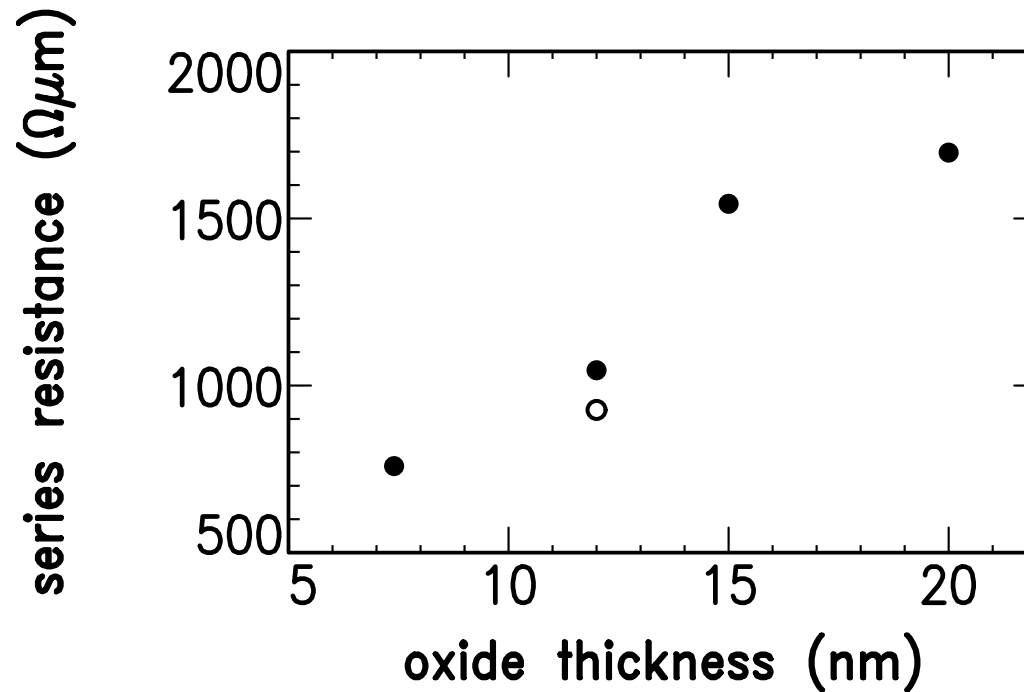


- process dependence of the carrier mobility for N-channels

$$\mu_o = \frac{\beta_{\square}}{C_{ox}}$$



- process dependence of the carrier mobility reduction parameter  $\theta_1$  for infinite long and wide N-channel transistor ( $\Rightarrow \theta_{10}$ )



- process dependence of the series resistance for N-channels

$$R_s + R_d = \frac{\rho}{W_{\text{eff}}} \Rightarrow \rho = \frac{S_{L; \theta_1, R}}{\beta_{\square}}$$

- inversion charge  $Q_i$  from Poisson equation

$$Q_i \approx -C_{ox} \left( V_{gs} - V_{FB} - \phi_B - k \sqrt{V_{sb} + \phi_B} \right) = -C_{ox} (V_{gs} - V_{T1})$$

- threshold voltage  $V_{T1}$

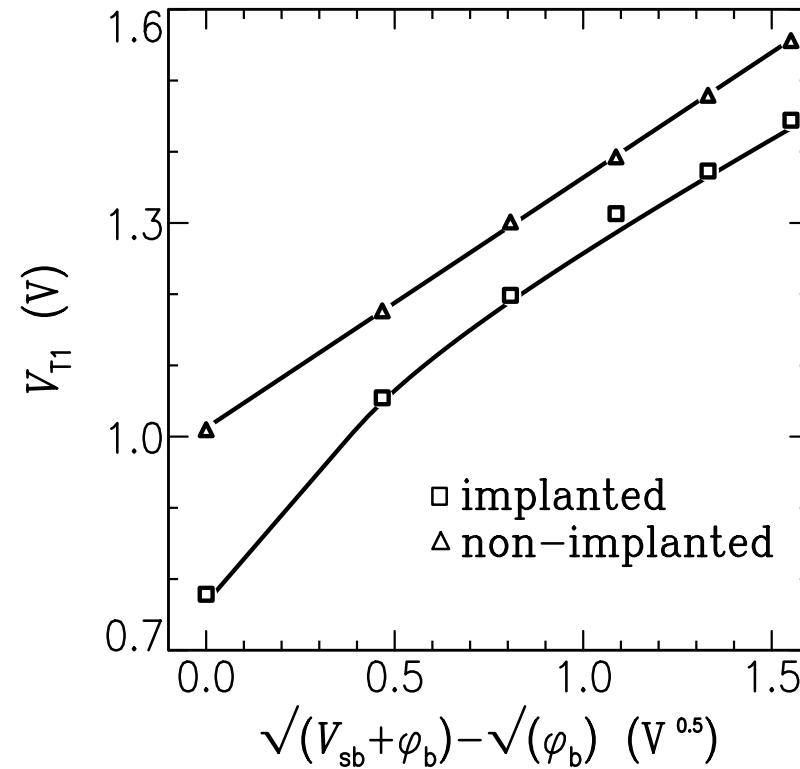
$$V_{T1} = V_{FB} + \phi_B + k \sqrt{V_{sb} + \phi_B} = V_{T0} + k \left( \sqrt{V_{sb} + \phi_B} - \sqrt{\phi_B} \right)$$

- MM9 threshold voltage parameter  $V_{T0}$

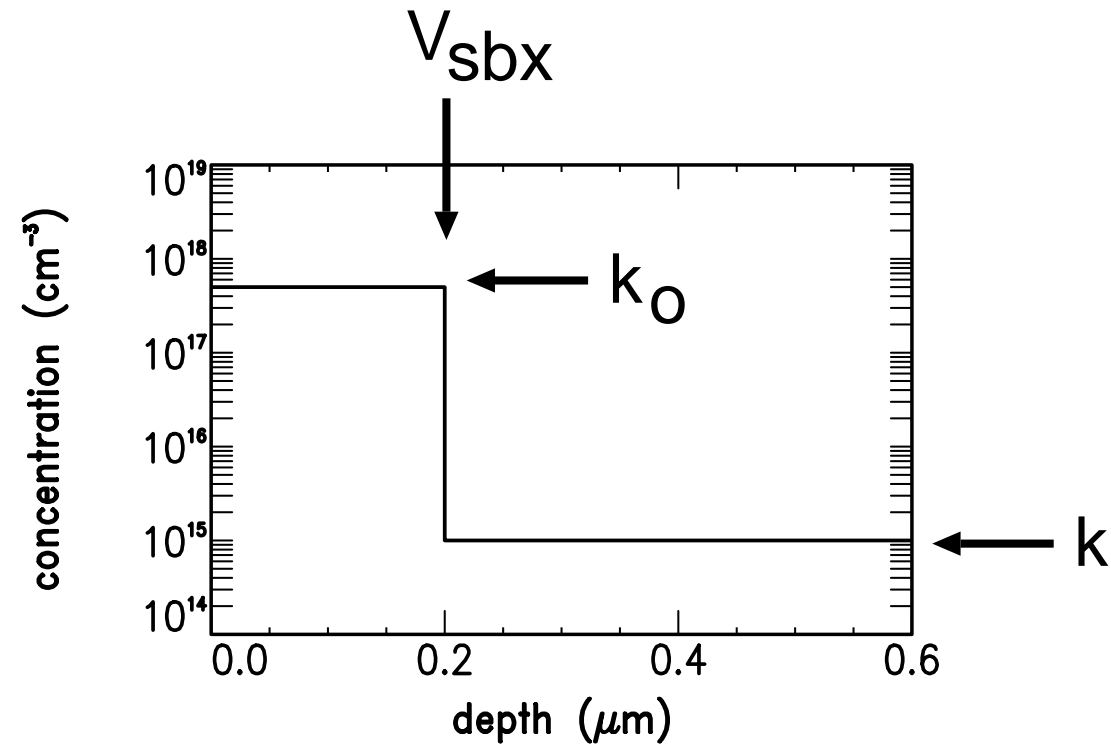
$$V_{T0} = V_{FB} + \phi_B + k \sqrt{\phi_B}$$

- $V_{FB} \Rightarrow$  flat band voltage  
 $\phi_B \Rightarrow$  surface potential in strong inversion  $\approx 2 \times$  bulk Fermi potential
- MM9 body-effect factor  $k$  depends on dopant concentration  $N$

$$k = \frac{\sqrt{2 \epsilon_s q N}}{C_{ox}}$$



- different behaviour for devices with a channel-implantation



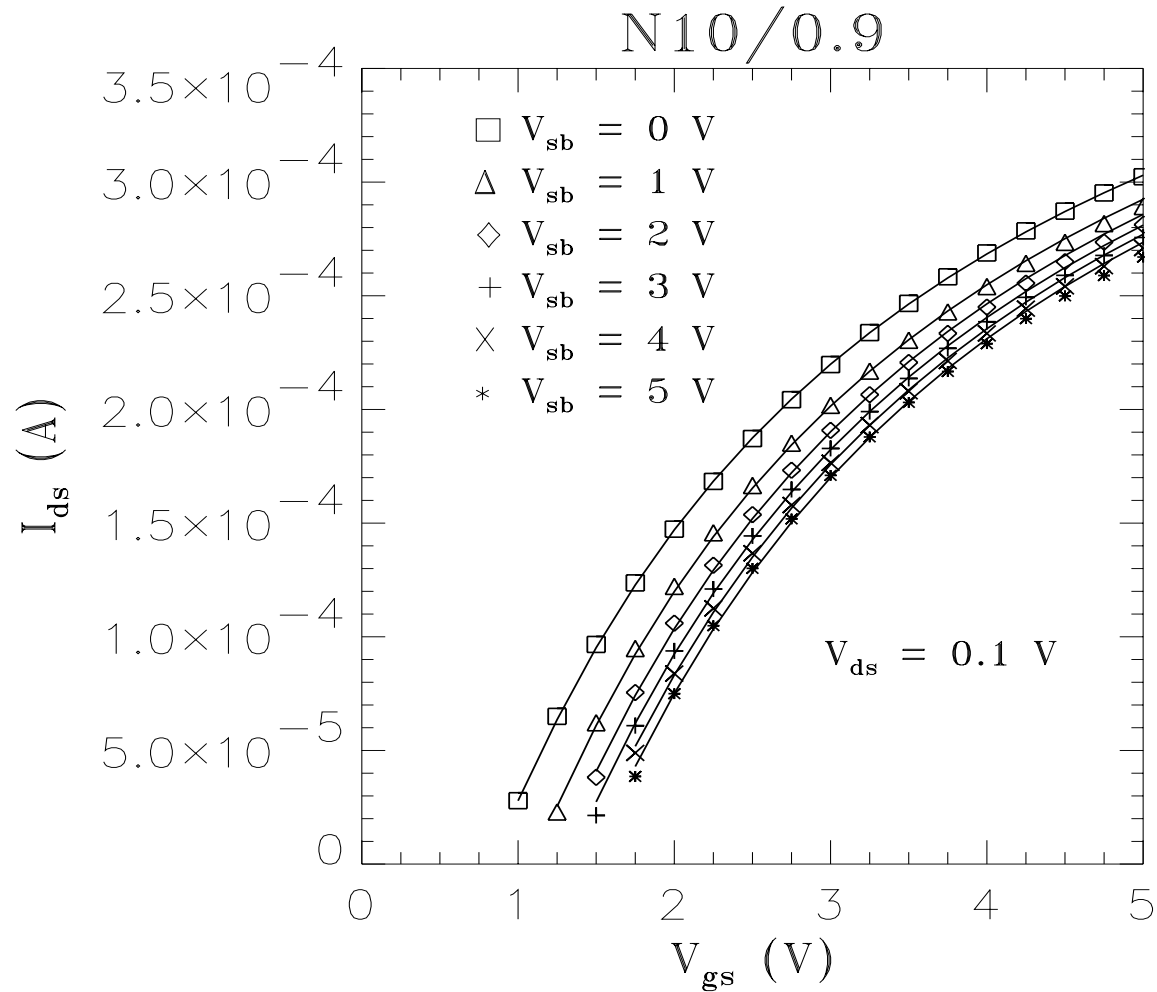
- with channel-implantation two MM9 body-effect factors:  $k_0$  and  $k$
- MM9 parameter  $V_{sbx}$  depends on depth of channel-implantation

- back-bias affects electrical field perpendicular to surface

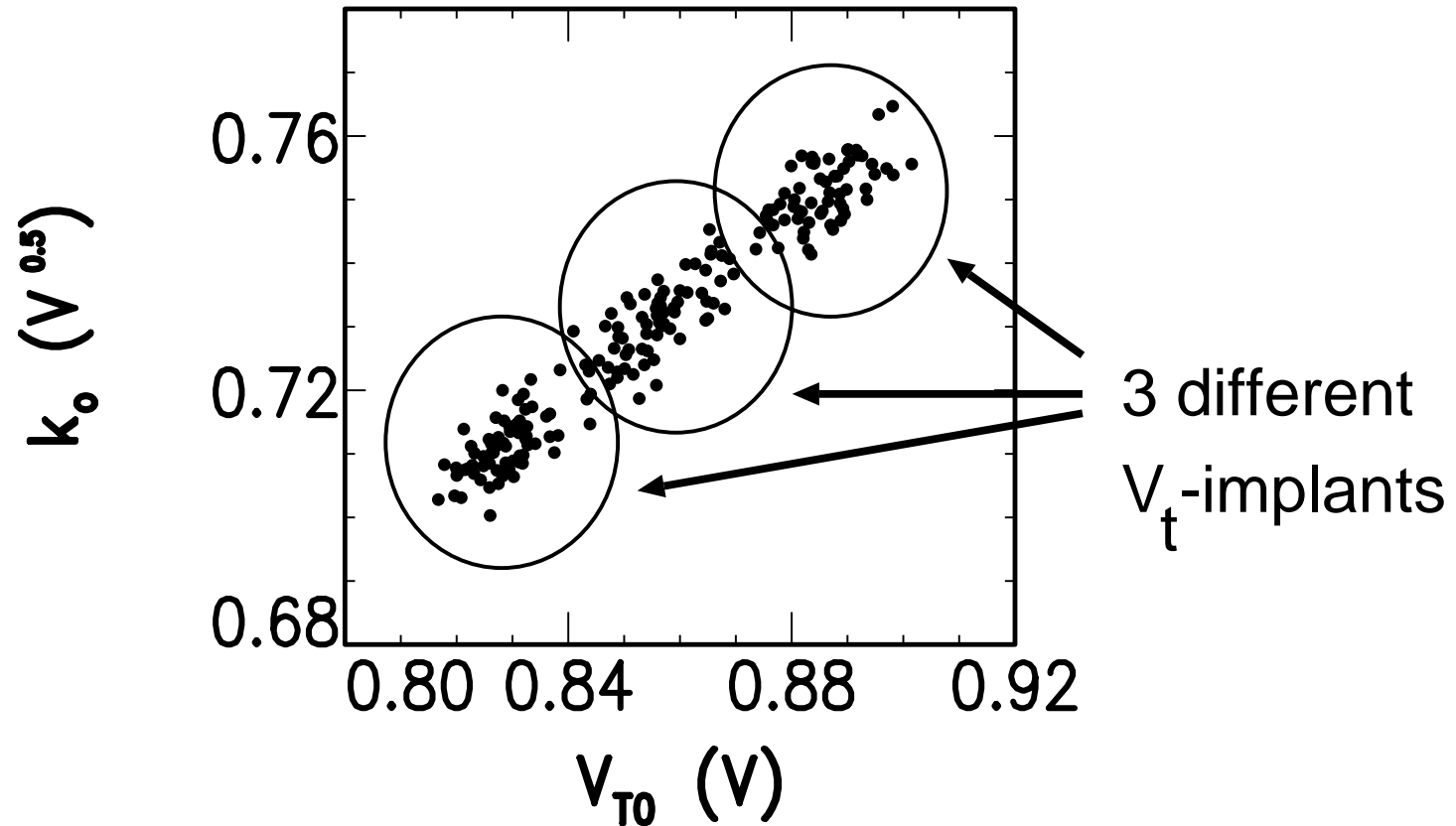
$$\mu = \frac{\mu_o}{1 + \theta_1 (V_{gs} - V_{T1}) + \theta_2 (\sqrt{V_{sb} + \phi_B} - \sqrt{\phi_B})}$$

- $\theta_2 \Rightarrow$  second MM9 mobility reduction parameter
- MM9 expression for  $I_{ds}$

$$I_{ds} = \beta \frac{(V_{gs} - V_{T1}) V_{ds} - \left(\frac{1 + \delta_1}{2}\right) V_{ds}^2}{1 + \theta_1 (V_{gs} - V_{T1}) + \theta_2 (\sqrt{V_{sb} + \phi_B} - \sqrt{\phi_B})}$$



MOS MODEL 9  
linear region



- both  $V_{T0}$  and  $k_0$  depend on  $V_T$ - or channel-implantation

- diffusion current  $I_{ds}$  from continuity equation

$$I_{ds} = -D W_{\text{eff}} \frac{\partial Q_i}{\partial x} = -\phi_T \mu W_{\text{eff}} \frac{\partial Q_i}{\partial x}$$

- integration over  $x$

$$\int_0^{L_{\text{eff}}} I_{ds} dx = L_{\text{eff}} I_{ds} = -\phi_T \mu W_{\text{eff}} \int_0^{L_{\text{eff}}} \frac{\partial Q_i}{\partial x} dx$$

$$L_{\text{eff}} I_{ds} = -\phi_T \mu W_{\text{eff}} \{Q_{iD} - Q_{iS}\}$$

- intermediate result

$$I_{ds} = \phi_T \mu \frac{W_{\text{eff}}}{L_{\text{eff}}} \{Q_{iS} - Q_{iD}\}$$

- $\mu \Rightarrow$  carrier mobility;  $Q_i \Rightarrow$  inversion charge per unit surface;  
 $\phi_T \Rightarrow$  thermal voltage

- intermediate result

$$I_{ds} = \phi_T \mu \frac{W_{eff}}{L_{eff}} \{Q_{iS} - Q_{iD}\}$$

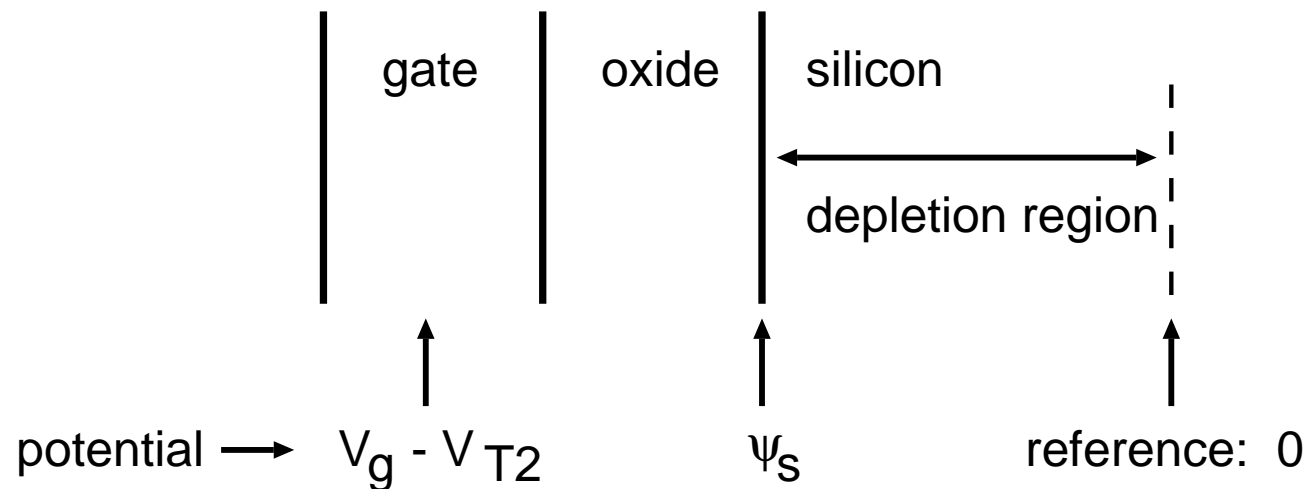
- inversion charge  $Q_i$  from Poisson equation

$$Q_i \approx Q_{i0} \exp\left(\frac{\psi_s}{\phi_T}\right) \Rightarrow Q_{iD} = Q_{iS} \exp\left(-\frac{V_{ds}}{\phi_T}\right)$$

- substitution

$$I_{ds} = \phi_T \mu \frac{W_{eff}}{L_{eff}} Q_{i0} \exp\left(\frac{\psi_{sS}}{\phi_T}\right) \left\{1 - \exp\left(-\frac{V_{ds}}{\phi_T}\right)\right\}$$

- $\psi_s \Rightarrow$  surface potential;  $\psi_{sS}$  surface potential at source-side



- dependence of surface potential on gate voltage

$$\psi_{sS} = \frac{C_{ox} (V_{gs} - V_{T2})}{C_{ox} + C_{depl}} = \frac{V_{gt2}}{m}$$

- the subthreshold slope  $m$  is given by

$$m = 1 + \frac{C_{depl}}{C_{ox}}$$

- from “theory”

$$I_{ds} = \phi_T \mu \frac{W_{eff}}{L_{eff}} Q_{i0} \exp\left(\frac{V_{gt2}}{m \phi_T}\right) \left\{ 1 - \exp\left(-\frac{V_{ds}}{\phi_T}\right) \right\}$$

- MM9 expression for  $I_{ds}$

$$I_{ds} = \beta G_3 \frac{V_{gt3} V_{ds1} - \left(\frac{1+\delta_1}{2}\right) V_{ds1}^2}{\left\{ 1 + \theta_1 V_{gt1} + \theta_2 (\sqrt{V_{sb} + \phi_B} - \sqrt{\phi_B}) \right\} (1 + \theta_3 V_{ds1})}$$

- MM9 expression for  $I_{ds}$

$$I_{ds} = \beta G_3 \frac{V_{gt3} V_{ds1} - \left(\frac{1+\delta_1}{2}\right) V_{ds1}^2}{\left\{1 + \theta_1 V_{gt1} + \theta_2 (\sqrt{V_{sb} + \phi_B} - \sqrt{\phi_B})\right\} (1 + \theta_3 V_{ds1})}$$

- herein  $V_{gt1} = \max\{0, (V_{gs} - V_{T1})\}$

$$G_1 = \exp\left(\frac{V_{gt2}}{2 m \phi_T}\right)$$

$$V_{gt3} = 2 m \phi_T \ln(1 + G_1)$$

$$V_{gt2} \ll -m \phi_T \Rightarrow V_{gt3} = 2 m \phi_T \exp\left(\frac{V_{gt2}}{2 m \phi_T}\right)$$

$$V_{gt2} \gg -m \phi_T \Rightarrow V_{gt3} = V_{gt2}$$

$$V_{ds1} = \frac{1}{1+\delta_1} V_{gt3}$$

$$G_3 = \frac{\zeta_1 \left\{1 - \exp\left(-\frac{V_{ds}}{\phi_T}\right)\right\} + G_1}{\frac{1}{\zeta_1} + G_1} \Rightarrow \zeta_1^2 \left\{1 - \exp\left(-\frac{V_{ds}}{\phi_T}\right)\right\}$$

- MM9 expression for  $I_{ds}$

$$I_{ds} = \zeta_1^2 \beta \left\{ 1 - \exp\left(-\frac{V_{ds}}{\phi_T}\right) \right\} \frac{\frac{1}{2} \frac{1}{1+\delta_1} V_{gt3}^2}{\{1 + \theta_2 (\sqrt{V_{sb} + \phi_B} - \sqrt{\phi_B})\}}$$

$$I_{ds} = \zeta_1^2 \beta \left\{ 1 - \exp\left(-\frac{V_{ds}}{\phi_T}\right) \right\} \frac{\frac{2 m^2 \phi_T^2}{1+\delta_1} \exp\left(\frac{V_{gt2}}{m \phi_T}\right)}{\{1 + \theta_2 (\sqrt{V_{sb} + \phi_B} - \sqrt{\phi_B})\}}$$

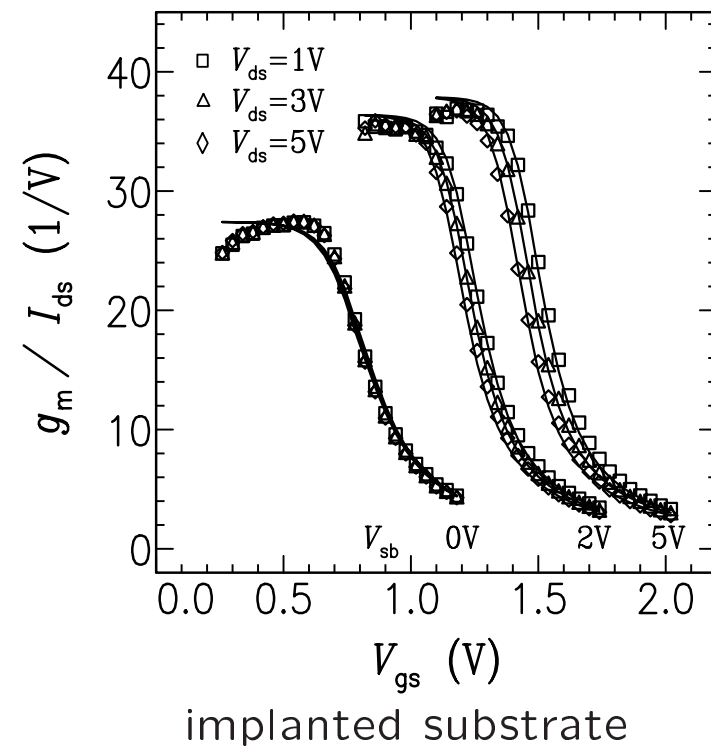
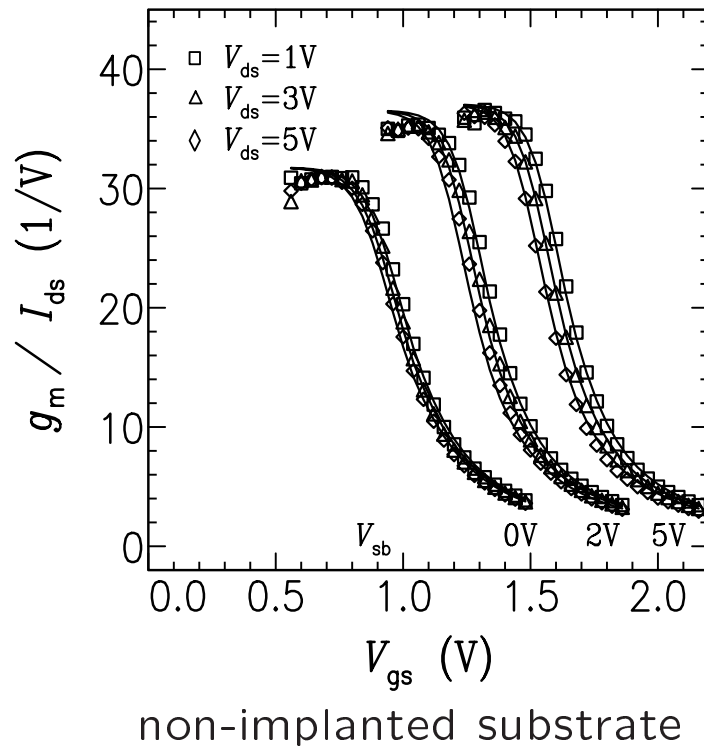
- from “theory”

$$I_{ds} = \phi_T \mu \frac{W_{eff}}{L_{eff}} Q_{i0} \exp\left(\frac{V_{gt2}}{m \phi_T}\right) \left\{ 1 - \exp\left(-\frac{V_{ds}}{\phi_T}\right) \right\}$$

- $\zeta_1$  is the “weak-inversion correction” parameter of MM9

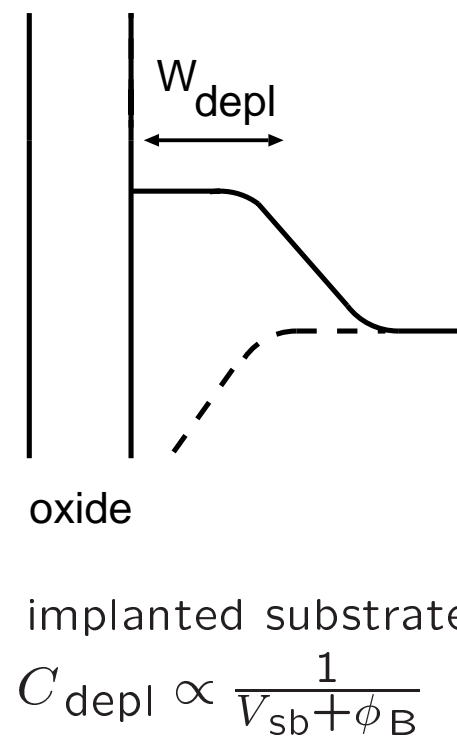
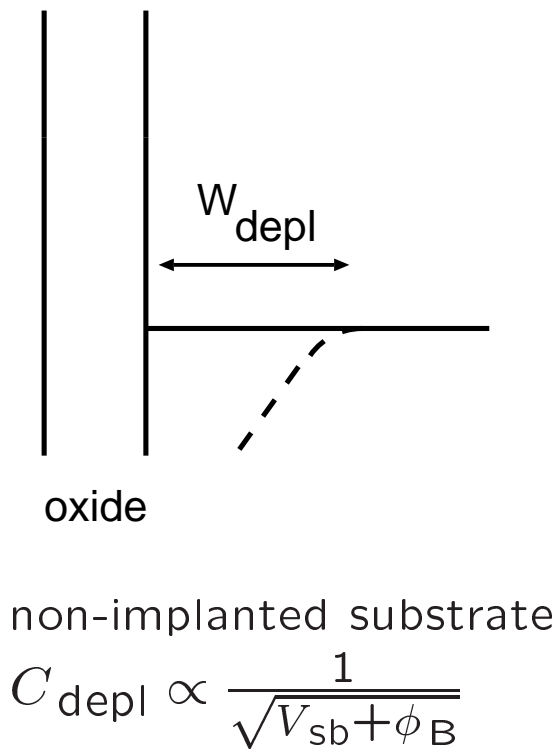
- the subthreshold slope  $m$  can be determined from

$$m = \left( \phi_T \frac{d \ln I_{ds}}{d V_{gs}} \right)^{-1} = \left( \phi_T \frac{g_m}{I_{ds}} \right)^{-1}$$



- the subthreshold slope  $m$  is given by

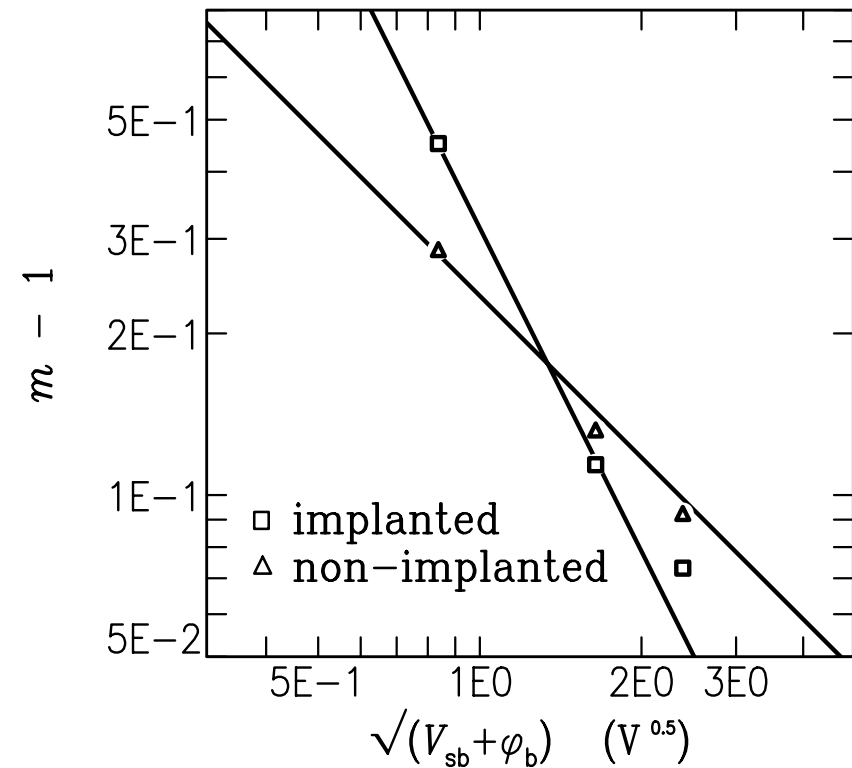
$$m = 1 + \frac{C_{\text{depl}}}{C_{\text{ox}}}$$



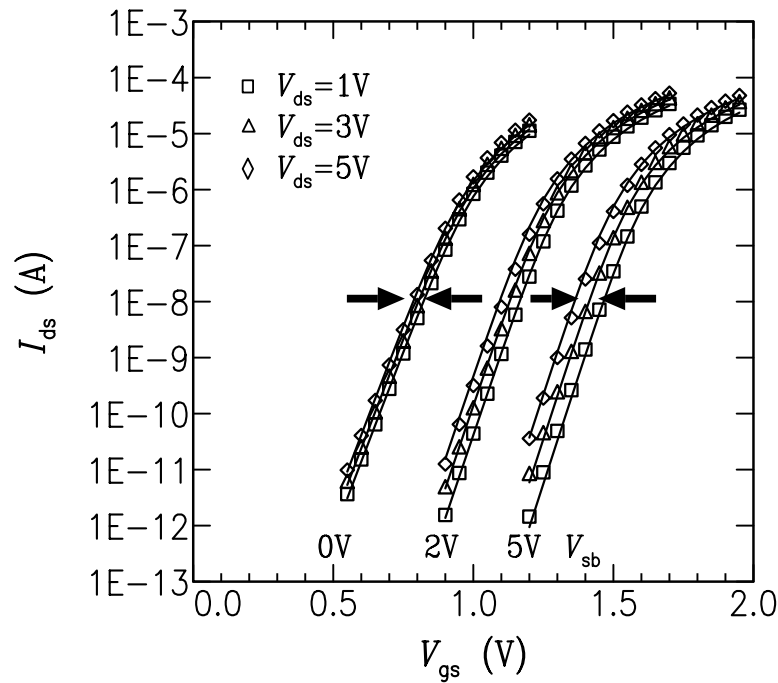
- in MM9 the subthreshold slope  $m$  is modelled by

$$m = 1 + m_0 \left( \frac{\sqrt{\phi_b}}{\sqrt{V_{sb} + \phi_b}} \right)^{\eta_m}$$

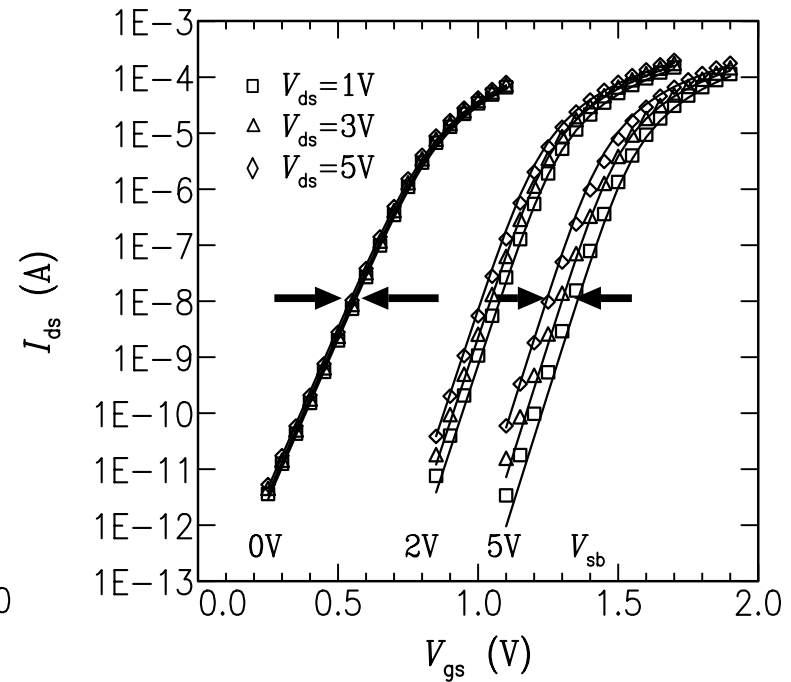
- non-impl. substrate  $\Rightarrow \eta_m = 1$
- implanted substrate  $\Rightarrow \eta_m = 2$
- $m_0 \Rightarrow$  MM9 subthreshold-slope parameter



- threshold shift increasing with  $V_{ds}$

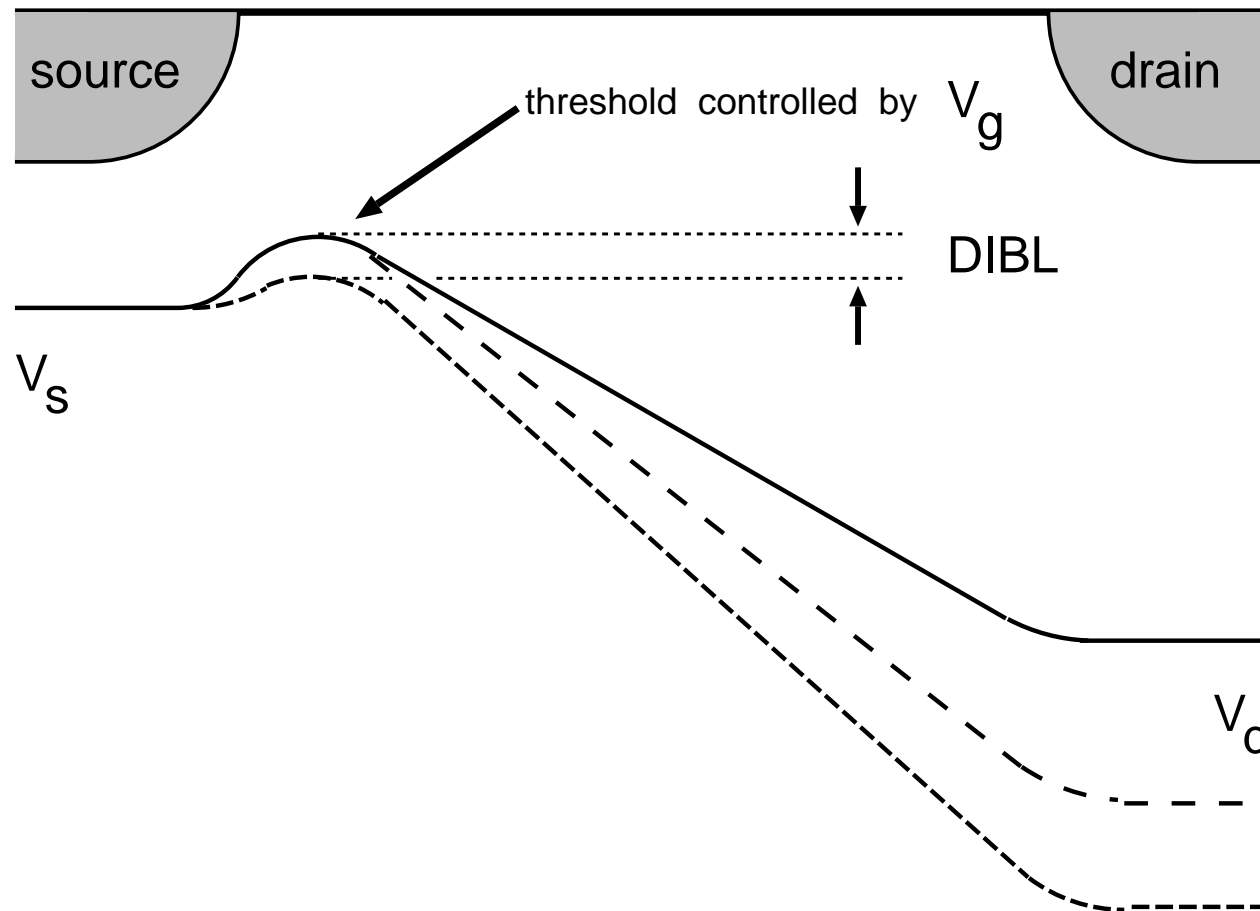


non-implanted substrate



implanted substrate

- threshold shift increasing with  $V_{ds}$



- theoretical drain-induced barrier-lowering for non-implanted substrate

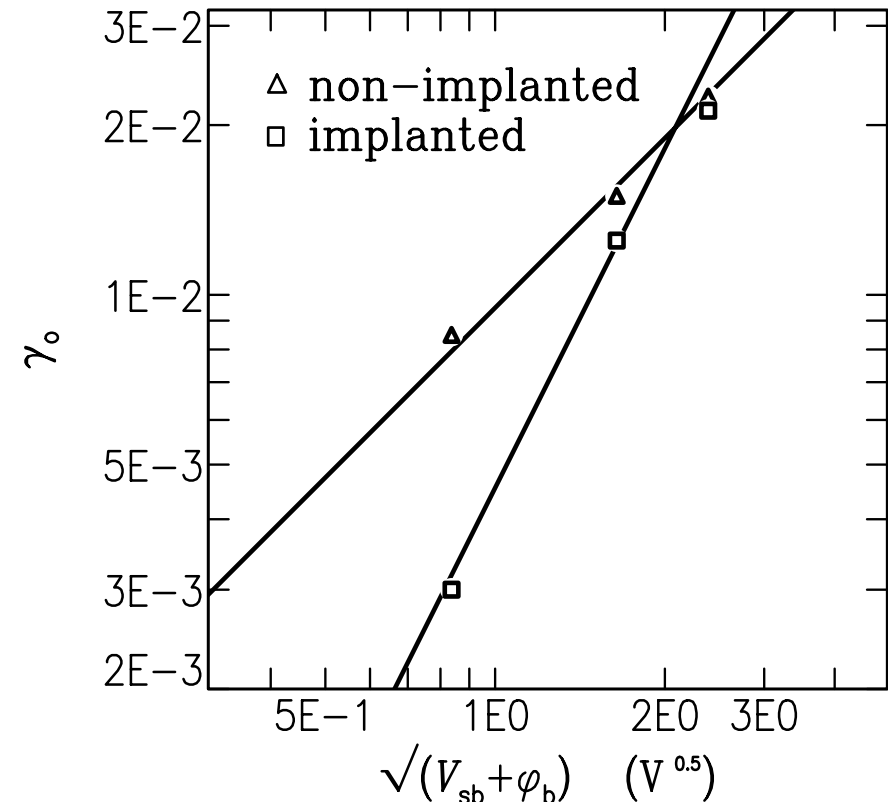
$$\Delta V_{T, \text{DIBL}} \propto - \frac{\sqrt{V_{sb} + \phi_b}}{L_{\text{eff}}^2} V_{ds}$$

- DIBL in MOS MODEL 9

$$\Delta V_{T, \text{DIBL}} = -\gamma_0 V_{ds}$$

$$\gamma_0 = \gamma_{00} \left( \frac{\sqrt{V_{sb} + \phi_b}}{\sqrt{\phi_b}} \right)^{\eta_\gamma}$$

- non-impl. substrate  $\Rightarrow \eta_\gamma = 1$
- implanted substrate  $\Rightarrow \eta_\gamma = 2$
- $\gamma_{00} \Rightarrow$  MM9 DIBL parameter



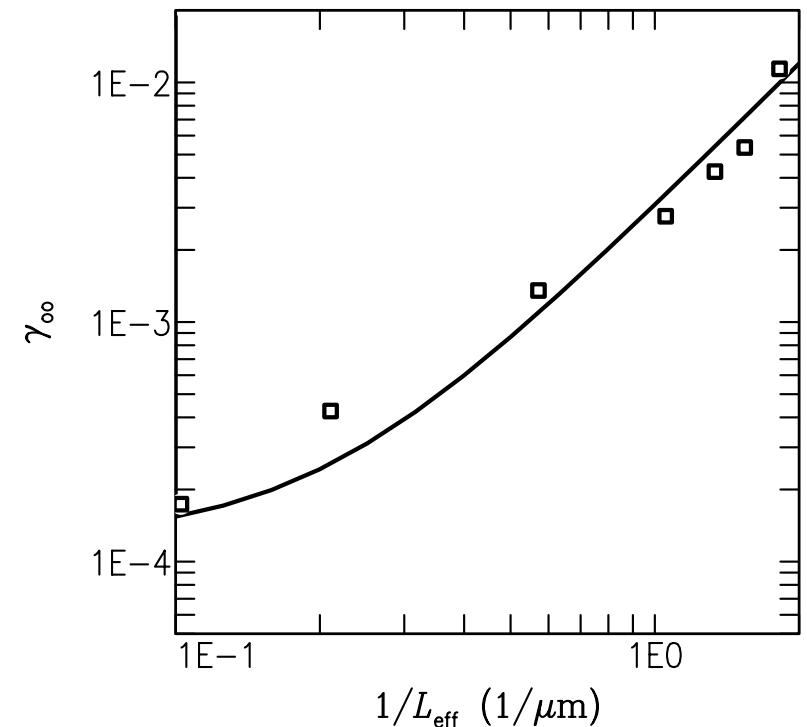
- theoretical drain-induced barrier-lowering for non-implanted substrate

$$\Delta V_{T, \text{DIBL}} \propto - \frac{\sqrt{V_{sb} + \phi_b}}{L_{\text{eff}}^2} V_{ds}$$

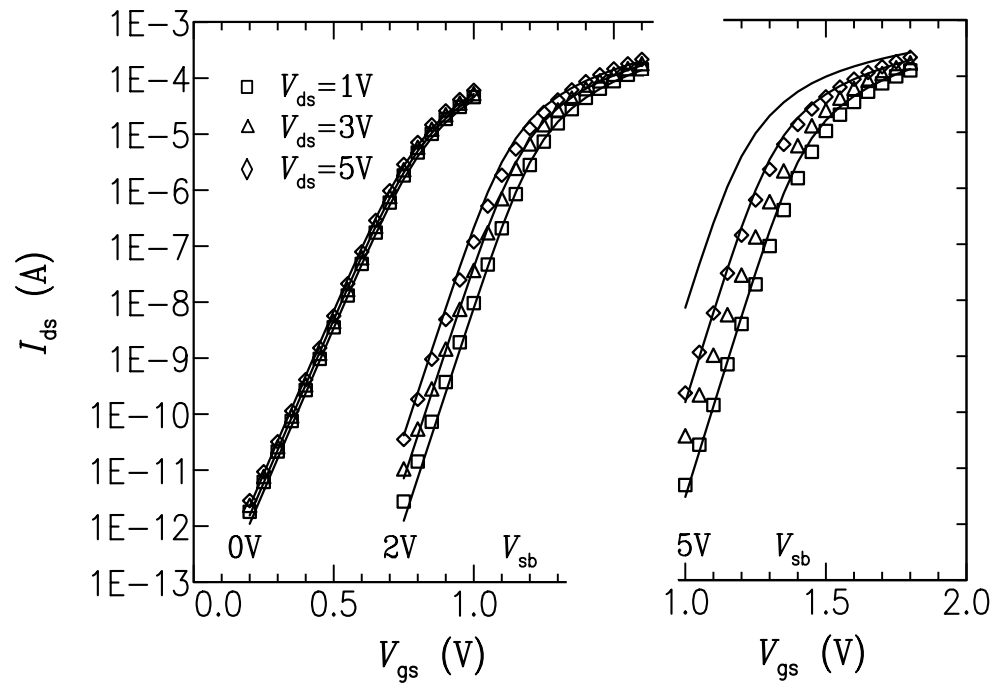
- scaling of DIBL in MM9

$$\gamma_{00} = \gamma_{00R} + \left( \frac{1}{L_{\text{eff}}^2} - \frac{1}{L_{\text{effR}}^2} \right) S L; \gamma_{00}$$

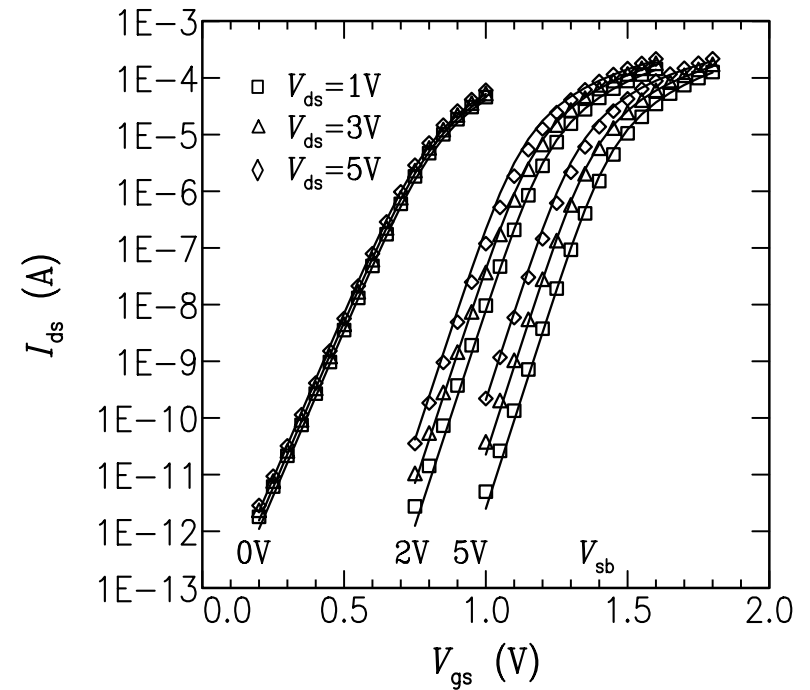
- example for implanted substrate



- very small channel-lengths: source & drain depletion regions almost touch

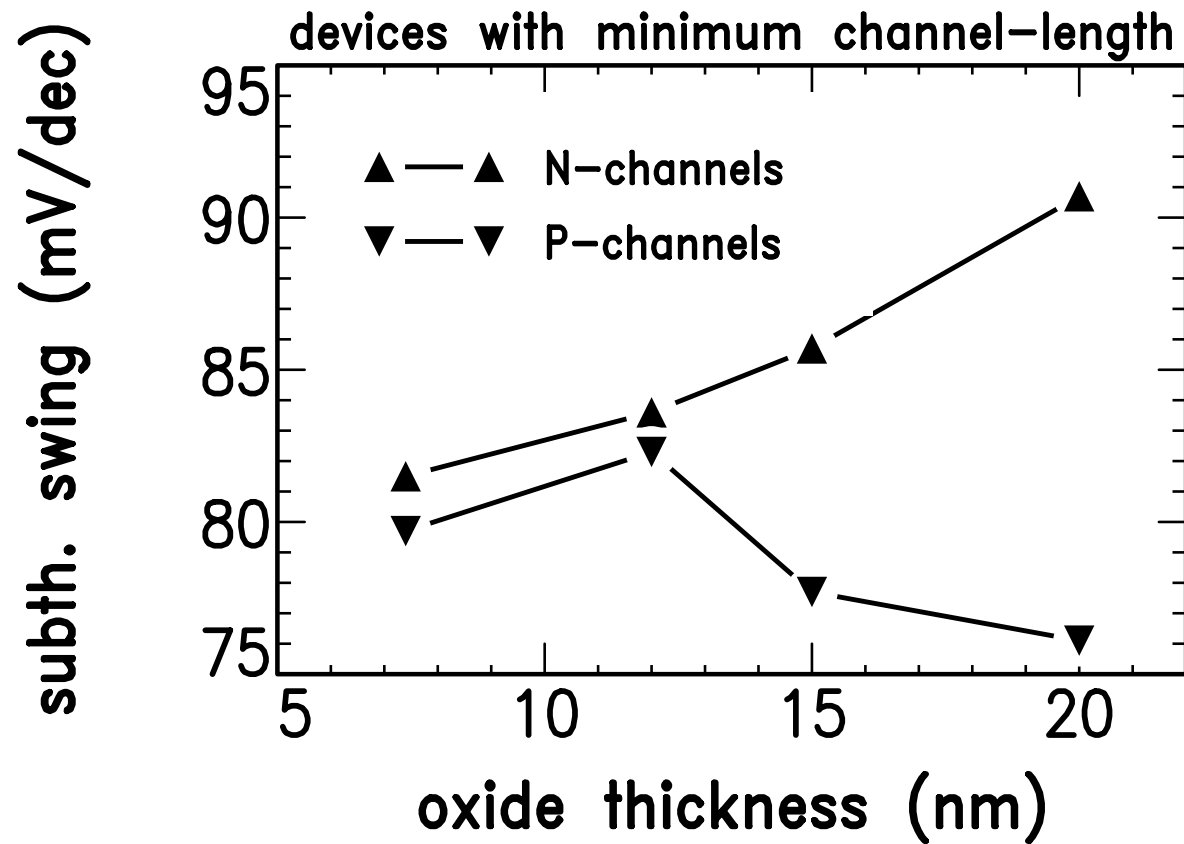


without upper limit  
for  $V_{sb}$  dependence

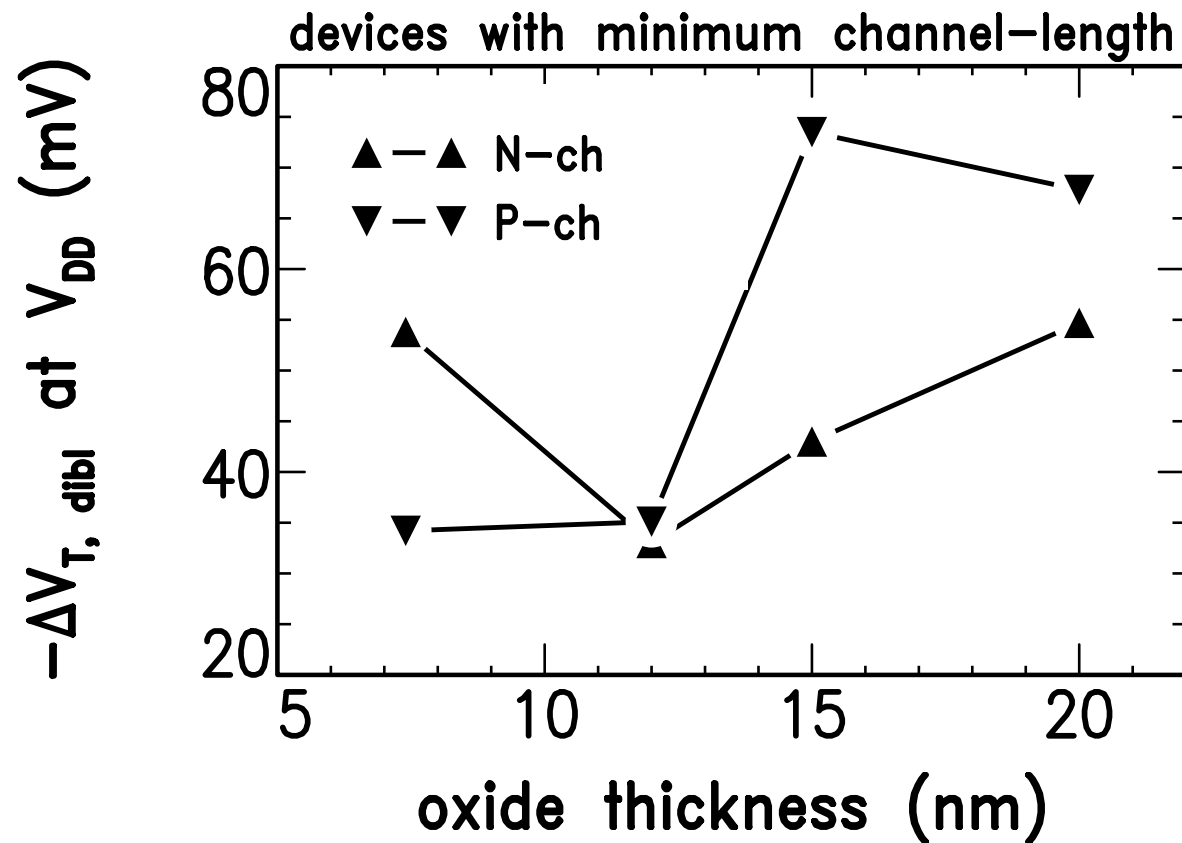


with upper limit  
for  $V_{sb}$  dependence

- $V_{sb} \leq V_{SBT}$ 
  - $\gamma_0 = \gamma_{00} \left( \frac{\sqrt{V_{sb} + \phi_b}}{\sqrt{\phi_b}} \right)^{\eta_\gamma}$
  - $m = 1 + m_0 \left( \frac{\sqrt{\phi_b}}{\sqrt{V_{sb} + \phi_b}} \right)^{\eta_m}$
- $V_{sb} > V_{SBT}$ 
  - $\gamma_0 = \gamma_{00} \left( \frac{\sqrt{V_{SBT} + \phi_b}}{\sqrt{\phi_b}} \right)^{\eta_\gamma}$
  - $m = 1 + m_0 \left( \frac{\sqrt{\phi_b}}{\sqrt{V_{SBT} + \phi_b}} \right)^{\eta_m}$
- non-impl. substrate  $\Rightarrow \eta_\gamma = \eta_m = 1$   
 implanted substrate  $\Rightarrow \eta_\gamma = \eta_m = 2$
- MM9 parameter  $V_{SBT}$  important only for smallest channel lengths

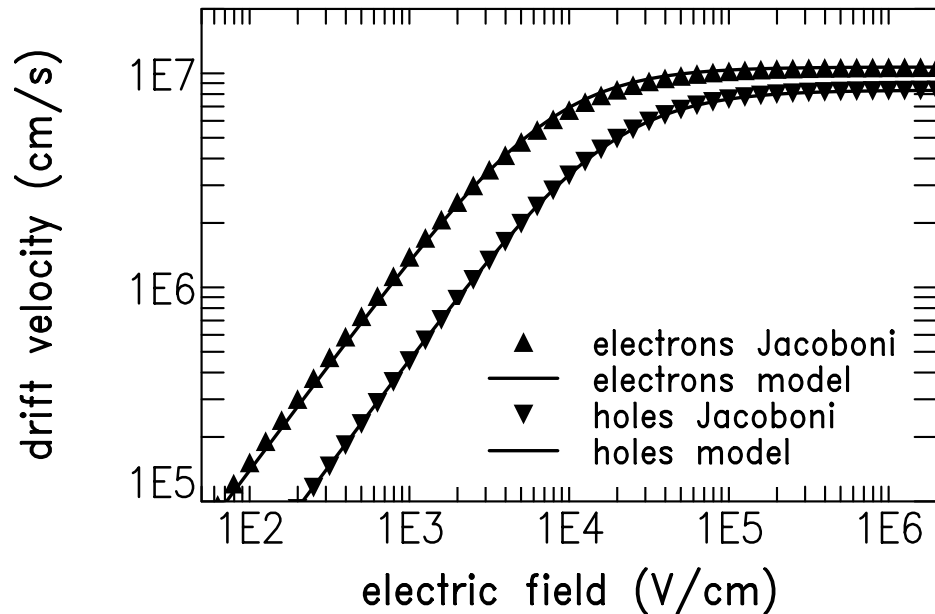


- subthreshold swing  $S = (1 + m_o) \phi_T \ln 10$



- drain-induced barrier-lowering at  $V_{dd}$ :  $\Delta V_{T, \text{DIBL}} = -\gamma_{00} V_{dd}$

- at high electrical fields the drift velocity saturates

drift velocity vs.  $E_{\parallel}$ 

- relation between drift velocity  $v$ , mobility  $\mu$  and electric field  $E_{\parallel}$

$$v = \mu(E_{\parallel}) E_{\parallel} = \frac{\mu E_{\parallel}}{1 + \frac{E_{\parallel}}{E_{\text{sat}}}} \Rightarrow \mu(E_{\parallel}) = \frac{\mu}{1 + \frac{E_{\parallel}}{E_{\text{sat}}}}$$

- carrier mobility depends on drain-source voltage

$$\mu(E_{\parallel}) = \frac{\mu}{1 + \frac{E_{\parallel}}{E_{\text{sat}}}} = \frac{\mu}{1 + \frac{V_{\text{ds}}}{L_{\text{eff}} E_{\text{sat}}}} = \frac{\mu}{1 + \theta_3 V_{\text{ds}}}$$

- $\theta_3$  is the third MM9 mobility reduction parameter

$$\theta_3 = \frac{1}{L_{\text{eff}} E_{\text{sat}}}$$

- drift current  $I_{\text{ds}}$

$$I_{\text{ds}} = -\mu W_{\text{eff}} Q_i \frac{\partial \phi}{\partial x} = -v W_{\text{eff}} Q_i$$

saturates due to pinch-off and velocity saturation

- expression for  $V_{ds, sat}$  from expression for  $I_{ds}$

$$I_{ds} = \mu C_{ox} \frac{W_{eff}}{L_{eff}} \left[ (V_{gs} - V_{TO}) V_d - \frac{1}{2} V_d^2 \right] \Rightarrow V_{ds, sat} = V_{gs} - V_{TO}$$

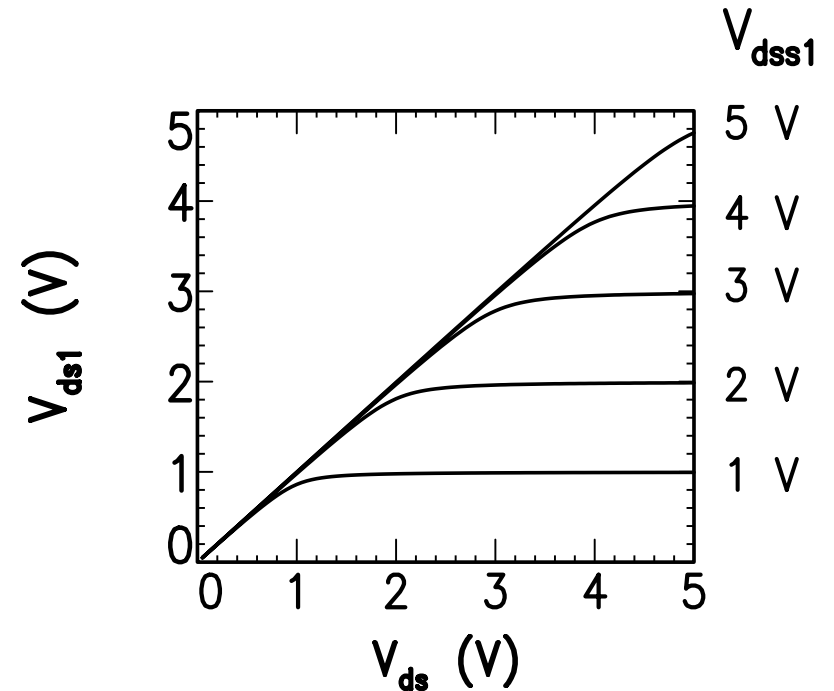
- MM9 expression for  $V_{ds, sat}$

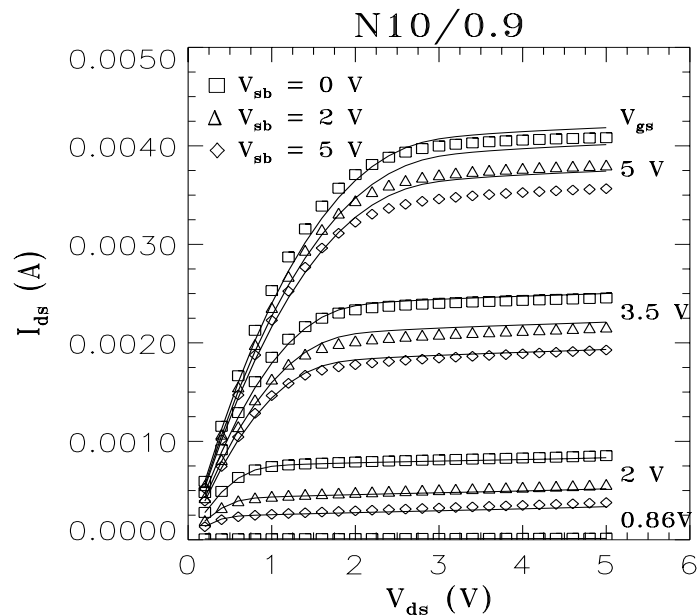
$$V_{dss1} = \frac{V_{gt3}}{1 + \delta_1} \frac{2}{1 + \sqrt{1 + \frac{2\theta_3 V_{gt3}}{1 + \delta_1}}}$$

- effective source-gate voltage  $V_{gt3}$

$$V_{gt3} = V_{gs} - V_{T2} \approx V_{gs} - V_{T1}$$

- $V_{ds1}$  changes smoothly  
from  $V_{ds}$  to  $V_{dss1}$





MOS MODEL 9  
saturation region

- MM9 expression for  $I_{ds}$

$$I_{ds} = \beta \frac{V_{gt3} V_{ds1} - \left(\frac{1 + \delta_1}{2}\right) V_{ds1}^2}{\left\{1 + \theta_1 (V_{gs} - V_{T1}) + \theta_2 (\sqrt{V_{sb} + \phi_B} - \sqrt{\phi_B})\right\} (1 + \theta_3 V_{ds1})}$$

- drain series resistance is included in  $\theta_3$

$$\theta_3 \approx \theta_{30} - \beta R_d - \frac{1}{2} \theta_{10}$$

- drain series resistance is included in  $\theta_1$  and  $\theta_3$

$$\theta_1 = \theta_{10} + \beta (R_s + R_d)$$

$$\theta_3 \approx \theta_{30} - \beta R_d - \frac{1}{2} \theta_{10}$$

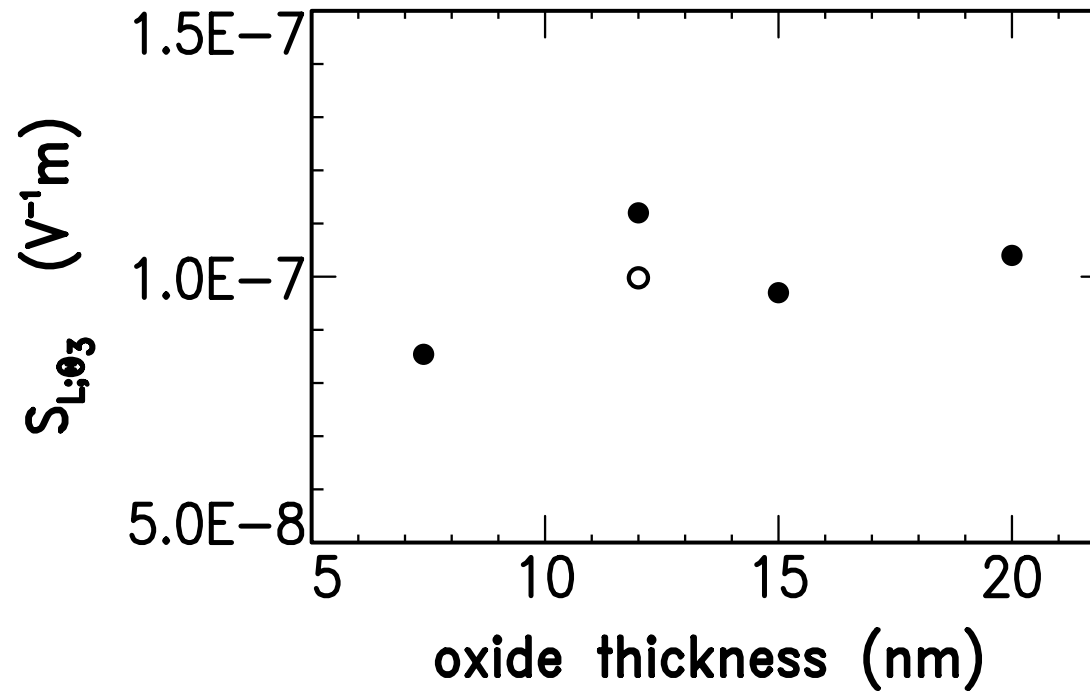
$$\theta_1 + \theta_3 \approx \frac{1}{2} \theta_{10} + \theta_{30} + \beta R_s$$

- MM9 expression for  $I_{ds}$  with  $V_{sb} = 0$

$$I_{ds} = \beta \frac{V_{gt3} V_{ds1} - \left(\frac{1+\delta_1}{2}\right) V_{ds1}^2}{\{1 + \theta_1 (V_{gs} - V_{T1})\} (1 + \theta_3 V_{ds1})}$$

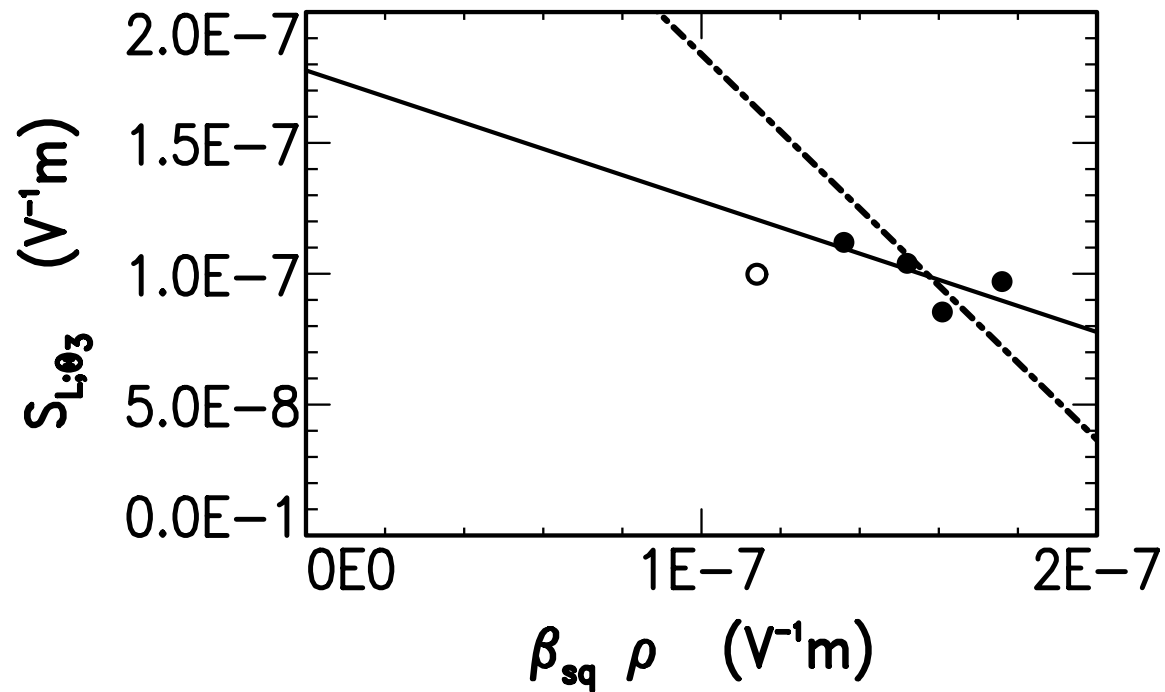
- with  $V_{ds1} = V_{dss1} \approx V_{gt3} \approx V_{gs} - V_{T1} \Rightarrow$  MM9 expression for  $I_{ds,sat}$

$$I_{ds,sat} \approx \beta \frac{\left(\frac{1-\delta_1}{2}\right) (V_{gs} - V_{T1})^2}{1 + (\theta_1 + \theta_3) (V_{gs} - V_{T1})}$$



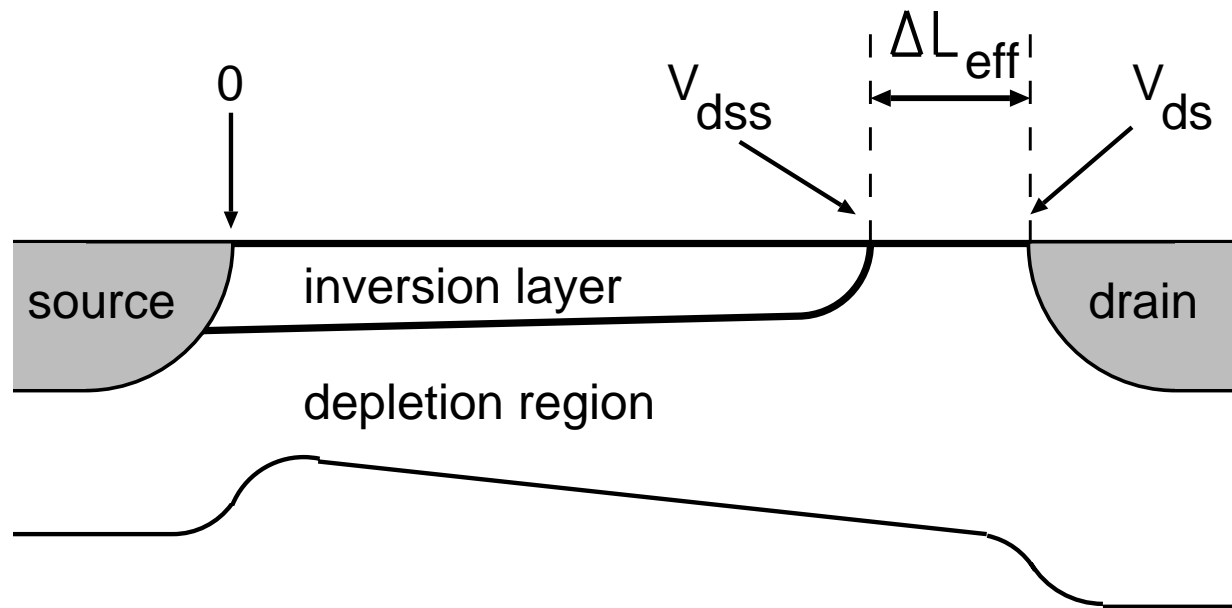
- process dependence of the length dependence of  $\theta_3$  for N-channels

$$\theta_3 \approx \theta_{30} - \beta R_d - \frac{1}{2} \theta_{10} \quad \Rightarrow \quad S_{L;\theta_3,R} \approx \frac{1}{E_{\text{sat}}} - \frac{\beta \square \rho}{2}$$



- process dependence of the length dependence of  $\theta_3$  for N-channels

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- at high drain-source voltages a further increase in drain-source voltage
  - reduces the channel length  $\Rightarrow$  channel-length modulation
  - increases the mobile charge in the channel  $\Rightarrow$  static feedback

- channel-length modulation  $\Delta L_{\text{eff}}$

$$\frac{\Delta L_{\text{eff}}}{L_{\text{eff}}} \approx \alpha \ln \left( 1 + \frac{V_{\text{ds}} - V_{\text{ds1}}}{V_{\text{p}}} \right)$$

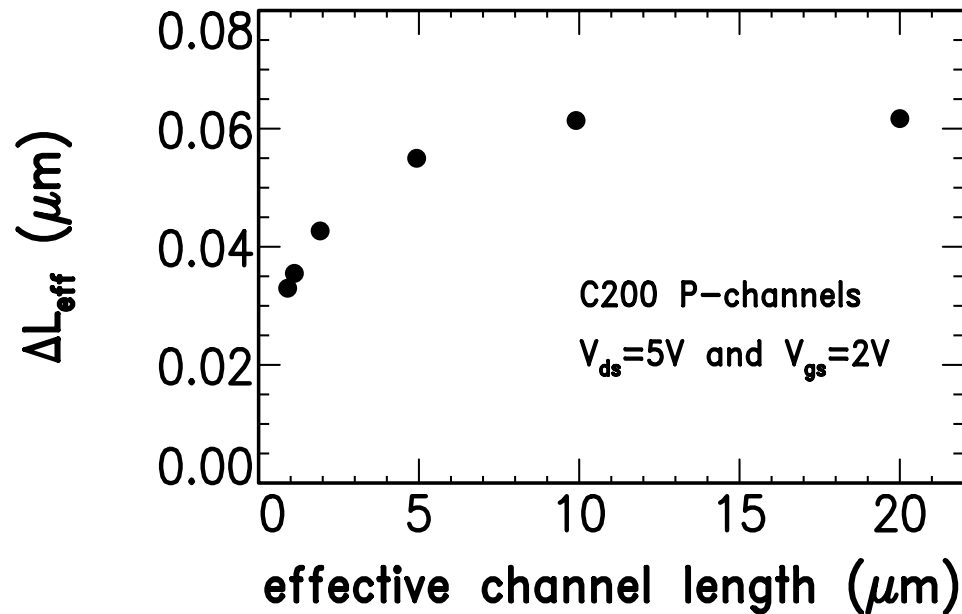
- channel-length modulation is modelled by multiplying drain current by

$$G_2 = 1 + \alpha \ln \left( 1 + \frac{V_{\text{ds}} - V_{\text{ds1}}}{V_{\text{p}}} \right)$$

- but only in strong inversion by incorporation in  $G_3$

$$G_3 = \frac{\zeta_1 \left\{ 1 - \exp \left( -\frac{V_{\text{ds}}}{\phi_{\text{T}}} \right) \right\} + G_1 G_2}{\frac{1}{\zeta_1} + G_1}$$

- $\alpha$  and  $V_{\text{p}}$  are the MM9 parameters for channel-length modulation



- for  $L_{\text{eff}} > 5 \mu\text{m}$   $\Delta L_{\text{eff}}$  is almost independent of channel length
- for  $L_{\text{eff}} < 5 \mu\text{m}$  static feedback is dominating!
- multiplication factor  $G_2$  is **not** independent of channel length!

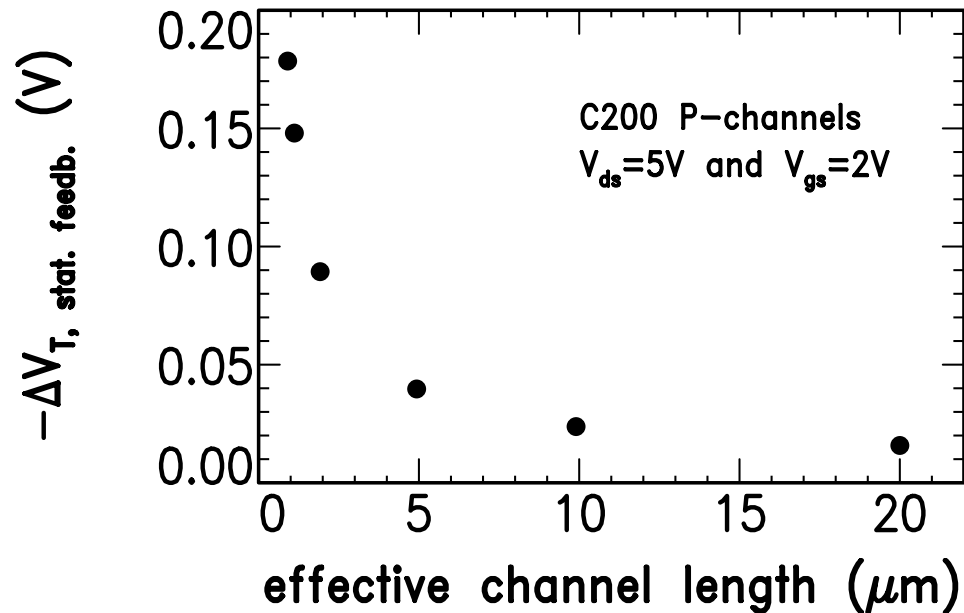
- static feedback is modelled by a shift in threshold voltage

$$\Delta V_{T, \text{stat. feedb.}} = -\gamma_1 V_{ds}^{0.6}$$

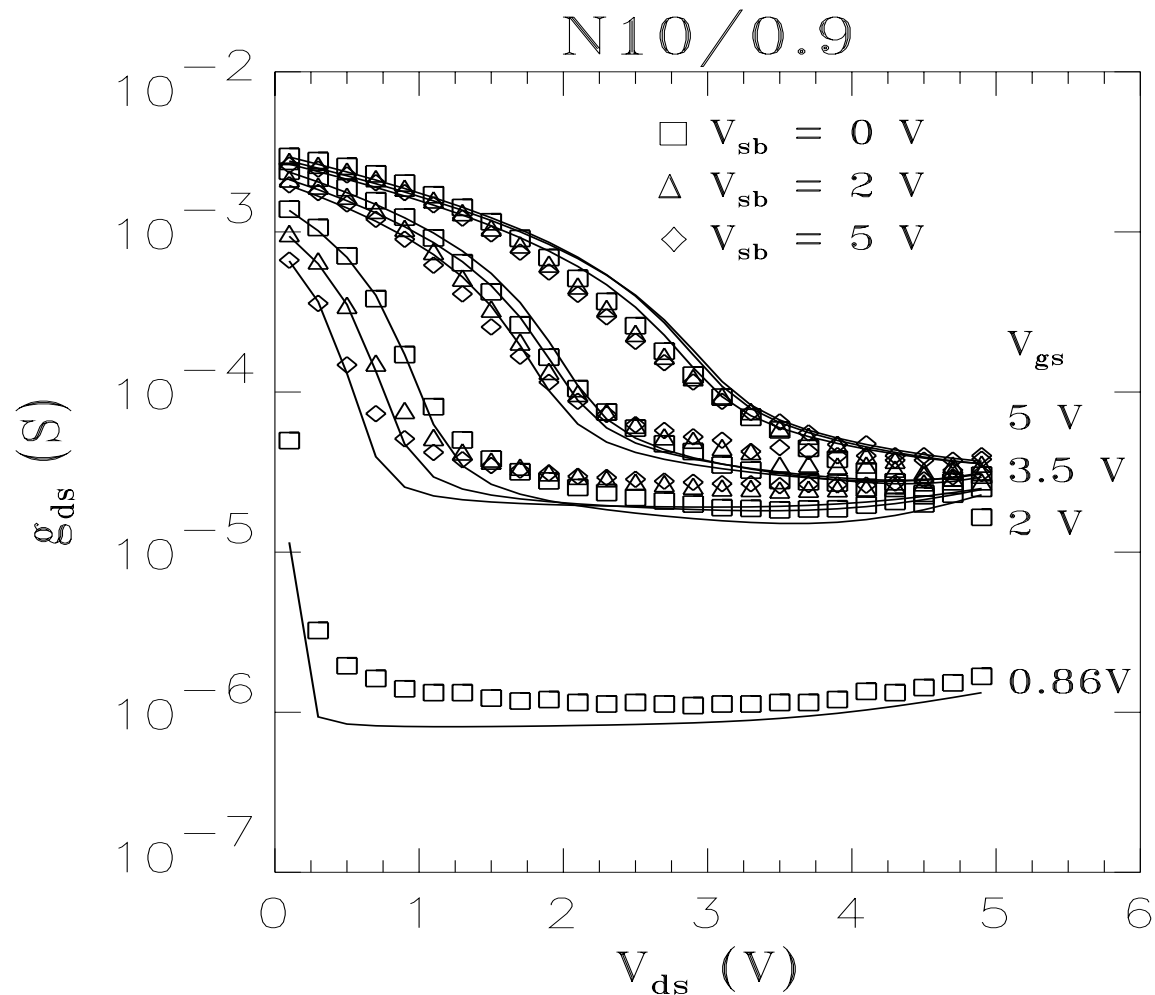
- static feedback is only effective in strong inversion

$$\Delta V_{T1} = \frac{V_{gtx}^2}{V_{gtx}^2 + V_{gt1}^2} \Delta V_{T, \text{dibl}} + \frac{V_{gt1}^2}{V_{gtx}^2 + V_{gt1}^2} \Delta V_{T, \text{stat. feedb.}}$$

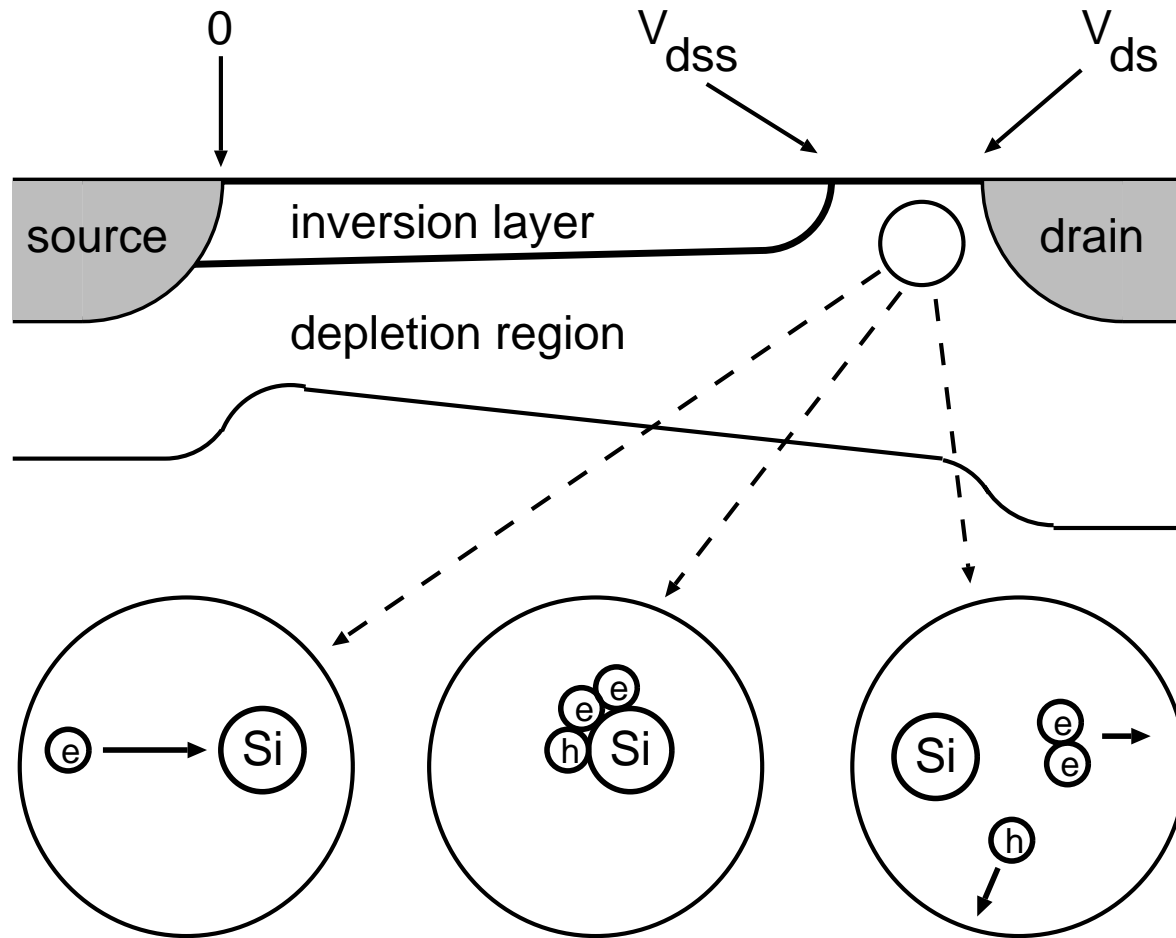
- herein  $V_{gt1} = \max\{0, (V_{gs} - V_{T1})\}$  and  $V_{gtx} = 0.707V$
- note:  $V_{T2} = V_{T1} + \Delta V_{T1}$
- $\gamma_1$  is the MM9 parameter for static feedback



- $\Delta V_{\text{stat. feedb.}}$  depends strongly on channel length
- for  $L_{\text{eff}} < 5 \mu\text{m}$  static feedback is dominating over channel-length modulation!



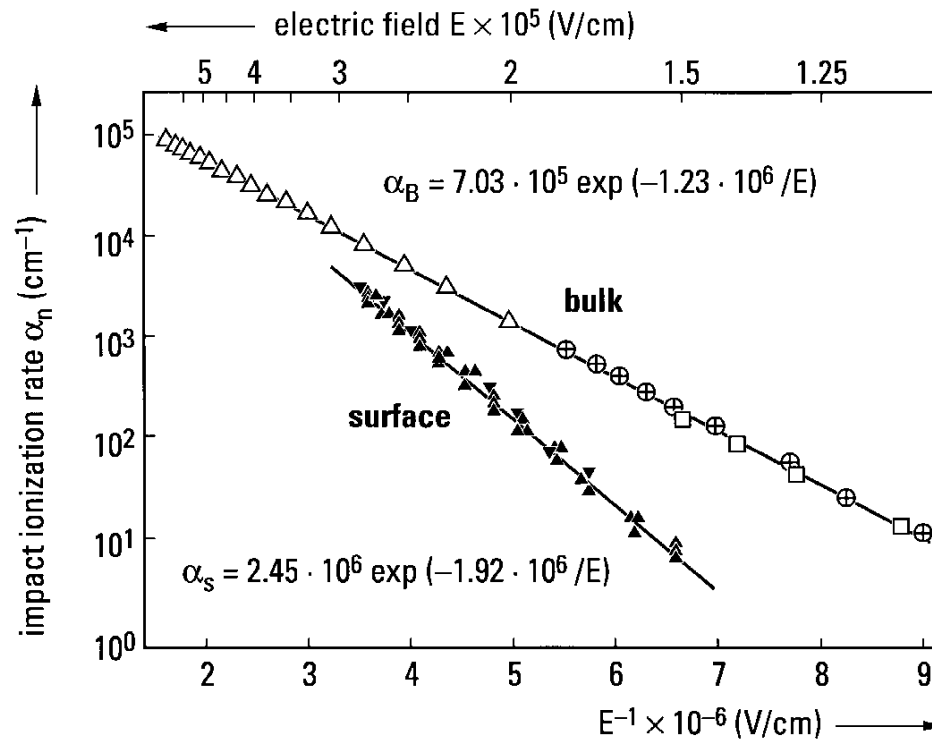
MOS MODEL 9  
output conductance



- substrate current is generated by impact-ionization

- generation by impact-ionization is given by

$$R_{ii} = \frac{-1}{q} (\alpha_n |J_n| + \alpha_p |J_p|)$$



impact-ionization rates  
for electrons vs.  $E_{\parallel}$

- impact-ionization varies strongly with electrical field

$$\alpha(E) = A \exp\left(-\frac{B}{E_{\parallel}}\right)$$

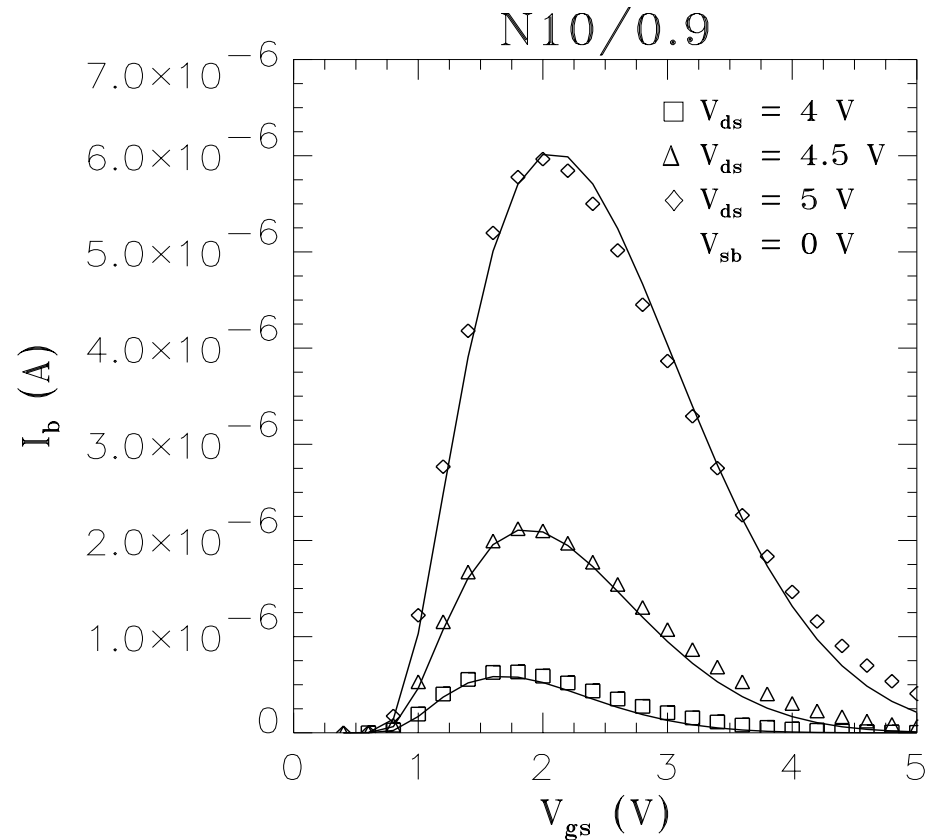
- weak avalanche:  
impact-ionization can be modelled as a multiplication factor

$$I_{avl} = a_1 I_{ds} \exp\left(\frac{-a_2}{V_{ds} - V_{dsa}}\right)$$

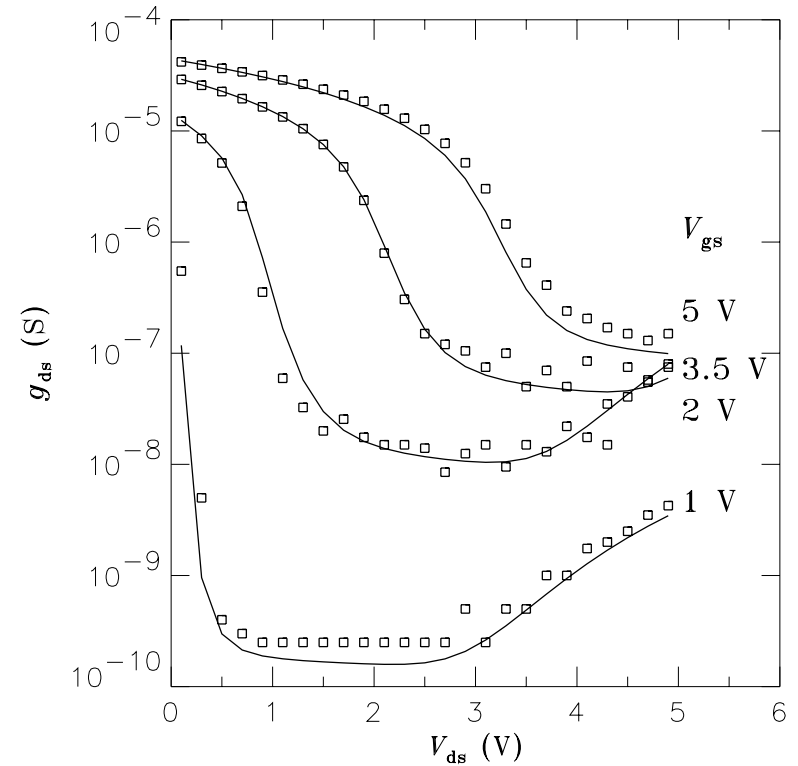
- sufficient high fields only in pinch-off region

$$V_{dsa} = a_3 V_{dss1} \quad \text{and} \quad I_{avl} = 0 \quad \text{if} \quad V_{ds} < V_{dsa}$$

- $I_{ds}$  and  $I_{avl}$  have to be added for total drain current
- $a_1$ ,  $a_2$  and  $a_3$  are the MM9 parameters for avalanche current



MOS MODEL 9 substrate current



N1.2/10: effect of substrate current on output conductance

## MOS MODEL 9

- relation between physics and equations
- no exact, detailed derivations
- trends in specific parameters
- description and no prediction of measurements

literature:

“Compact Transistor Modelling for Circuit Design”  
by H.C. de Graaff and F.M. Klaassen,  
Springer-Verlag, Wien-New York, 1990